Construction of two-dimensional quantum field models through Longo-Witten endomorphisms

Yoh Tanimoto

University of Rome "Tor Vergata", supported by Programma Rita Levi Montalcini

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Constructing QFT

What is quantum field theory? (cf. quantum mechanics, classical field)

- Lagrangian approach: perturbation theory is divergent.
- Need to compute *n*-point (Wightman) functions \implies S-matrix.
- Hilbert space and local observables can be reconstructed.

Form factor programme (Babujian, Karowski, Smirnov...)

- (Factorizing) S-matrix is conjectured from general requirements.
- Form factors $\operatorname{out}\langle q_1,\cdots,q_m|O(x)|p_1,\cdots,p_n
 angle^{\operatorname{in}}$ are obtained.
- *n*-point functions $\langle 0|O(x)O(0)|0\rangle = \sum_n \int dp_1 \cdots dp_n \langle 0|O(x)|p_1, \cdots, p_n \rangle^{\text{in in}} \langle p_1, \cdots, p_n | O(0)|0 \rangle.$
- Convergence? Locality in e.g. the sine-Gordon model?

von Neumann algebraic approach (Schroer, Lechner, T,...)

Find simpler observables in wedges, then local observables abstractly!

08/06/2017, Prague 2 / 10

Algebraic QFT

Haag-Kastler axioms

- Concerned with algebras of observables $\mathcal{A}(O)$ in spacetime regions O.
- Isotony, locality, Poincaré covariance, positivity of energy, existence of vacuum
- ϕ : quantum (Wightman) field $\Longrightarrow \mathcal{A}(O) := \overline{\{e^{i\phi(f)} : \operatorname{supp} f \subset O\}}^{\operatorname{vN}}$
- Examples: P(φ)₂, Yukawa₂, φ⁴₃, · · · (many other models in the Wightman axioms, Haag-Kastler net? c.f. Glimm-Jaffe)

Interacting QFT is difficult because pointlike field $\phi(x)$ is complicated... Isotony: $O_1 \subset O_2 \Longrightarrow \mathcal{A}(O_1) \subset \mathcal{A}(O_2)$ means that **larger** regions contain **more** observables, also **simpler** ones. Wedge: $W_R := \{(t, x) : x > |t|\}$.

A strategy for constructing Haag-Kastler nets

- \bullet Construct observables in ${\it W}_{\rm R}$
- Local observables by $\mathcal{A}(D_{a,b}) = \mathcal{A}(W_{\mathrm{R}} + a) \cap \mathcal{A}(W_{\mathrm{L}} + b)$

Standard wedge and double cone



Wedge-observables in integrable models (Schroer, Lechner)

- analytic S-matrix (e.g. the sinh-Gordon model) $S : \mathbb{R} + i(0, \pi) \to \mathbb{C}$, $\overline{S(\theta)} = S(\theta)^{-1} = S(-\theta) = S(\theta + \pi i), \ \theta \in \mathbb{R}.$
- S-symmetric Fock space: $\mathcal{H}_1 = L^2(\mathbb{R}, d\theta)$, $\mathcal{H}_n = P_n \mathcal{H}_1^{\otimes n}$, where P_n is the projection onto S-symmetric functions: $\Psi_n(\theta_1, \dots, \theta_n) = S(\theta_{k+1} - \theta_k)\Psi_n(\theta_1, \dots, \theta_{k+1}, \theta_k, \dots, \theta_n).$
- Zamolodchikov-Faddeev algebra: S-symmetrized creation and annihilation operators z[†](ξ) = Pa[†](ξ)P, z(ξ) = Pa(ξ)P, P = ⊕_n P_n.
- Wedge-local field: with J_1 : CPT, $f^{\pm}(\theta) = \int dx e^{\pm ix \cdot p(\theta)} f(x)$,

$$\phi(f) = z^{\dagger}(f^+) + z(J_1f^-).$$

Wedge-localization (Lechner '03, Bostelmann-Cadamuro '15)

If $\operatorname{supp} f$, $\operatorname{supp} g \subset W_{\mathrm{R}}$, then $[e^{i\phi(f)}, e^{iJ\phi(g)J}] = 0$. $\phi(f)$ formally commutes with observables coming from form factors localized at 0.

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Local observables

- Define $\mathcal{A}(W_{\mathrm{R}} + a) := \overline{\{e^{i\phi(f)} : \operatorname{supp} f \subset W_{\mathrm{R}} + a\}}^{\mathrm{vN}}$, $\mathcal{A}(W_{\mathrm{L}} + b) := \overline{\{e^{iJ\phi(f)J} : \operatorname{supp} f \subset W_{\mathrm{R}} - b\}}^{\mathrm{vN}}$.
- Question: is there any nontrivial observable in $\mathcal{A}(D_{a,b}) = \mathcal{A}(W_{\mathrm{R}} + a) \cap \mathcal{A}(W_{\mathrm{L}} + b)$?
- Let $\Delta^{\frac{1}{4}} = U(\Lambda(\frac{\pi i}{2}))$ (imaginary Lorentz boost), consider the map

$$\mathcal{A}(W_{\mathrm{R}}) \ni x \longmapsto \Delta^{\frac{1}{4}} U(a) x \Omega,$$

 $a \in W_{\rm R}$. If this is a **nuclear** map (approximated "well" by finite matrices), then there are **many** local observables (Buchhoz-D'Antoni-Longo '90).

Theorem (Lechner '08)

If S is analytic, satisfies a regularity condition and and S(0) = -1, the map above is indeed nuclear for sufficiently large a, therefore, there are local observables in $\mathcal{A}(D_{a,0})$ for such a, and hence a **Haag-Kastler net**.

S-matrices with poles (bound states)

If S has a pole (e.g. the Bullough-Dodd model, the sine-Gordon model), $\phi(f) = z^{\dagger}(f^+) + z(J_1f^-)$ is **no longer wedge-local**.

S: scalar, poles at $\theta = \frac{\pi i}{3}, \frac{2\pi i}{3}, S(\theta) = S\left(\theta + \frac{\pi i}{3}\right)S\left(\theta - \frac{\pi i}{3}\right)$ P_n : S-symmetrization, $\mathcal{H} = \bigoplus P_n \mathcal{H}_1^{\otimes n}, \mathcal{H}_1 = L^2(\mathbb{R}),$

$$(\chi_1(f))\xi(\theta) := \sqrt{2\pi|R|}f^+\left(\theta + \frac{\pi i}{3}\right)\xi\left(\theta - \frac{\pi i}{3}\right),$$

$$\chi_n(f) := n P_n(\chi_1(f) \otimes I \otimes \cdots \otimes I) P_n, \quad \chi(f) := \bigoplus \chi_n(f).$$

Theorem (Cadamuro-T. arXiv:1502.01313, CMP)

 $\widetilde{\phi}(f) := \phi(f) + \chi(f)$ commute weakly with $\widetilde{\phi}'(g)$ on a dense domain.

Problem: self-adjointness (problem of domain)! (\implies Haag-Kastler net) Work in progress: A_n -affine Toda, sine-Gordon...

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Twisting the massive free field

Let ϕ be the massive **complex** free field, with the charge operator Q, and $\mathcal{M} = \overline{\{e^{i\phi(f)} : \operatorname{supp} f \subset W_{\mathrm{R}}\}}^{\mathrm{vN}}$ be the wedge-algebra. Take $\mathcal{H} \otimes \mathcal{H}$.

Theorem (T. arXiv:1301.6090, FOMS)

From the algebra generated by $\mathcal{M} \otimes \mathbb{C}\mathbb{1}$ and $e^{tQ \otimes Q}(\mathbb{C}\mathbb{1} \otimes \mathcal{M})e^{-tQ \otimes Q}$, one can obtain an interacting Haag-Kastler net, $t \in \mathbb{R} \setminus 2\pi\mathbb{Z}$. Federbush?

Proof of wedge-localization: $e^{sQ} \cdot e^{-sQ}$ is an **automorphism** of \mathcal{M} . **A variation**: Let ϕ be the massive **real** free field, φ an inner symmetric function, \mathcal{M} the wedge-algebra. One can construct an operator \tilde{R}_{φ} :

Theorem (T., Alazzawi-Lechner '17)

From the algebra generated by $\mathcal{M} \otimes \mathbb{C}\mathbb{1}$ and $\tilde{R}_{\varphi}(\mathbb{C}\mathbb{1} \otimes \mathcal{M})\tilde{R}_{\varphi}^*$, one can obtain an interacting Haag-Kastler net.

Proof of wedge-localization: $\Gamma(\varphi(P_1)) \cdot \Gamma(\varphi(P_1))^*$ implements an endomorphism of \mathcal{M} . $\tilde{R}_{\varphi} \cdot \tilde{R}_{\varphi}^*$ maps $\mathcal{M} \otimes \mathbb{C}\mathbb{1}$ into $\mathcal{M} \otimes \mathcal{B}(\mathcal{H})$.

Relations to CFT?

 ϕ : massive free field, \mathcal{M} : the wedge algebra.

- One obtaines a CFT (the Heisenberg algebra) by restricting to the lightray.
- $\Gamma(\varphi(P_1))$ implements a Longo-Witten endomorphism: $\Gamma(\varphi(P_1))\mathcal{M}\Gamma(\varphi(P_1))^* \subset \mathcal{M} \text{ and } \Gamma(\varphi(P_1))$ commutes with the lightlike translations.

Take an interacting (integrable) QFT. Consider the wedge algebra \mathcal{M} and the restriction to the lightray $a \in W_{\rm B}$. **Question**: How large is the lightlike intersection $\mathcal{M} \cap U(a)\mathcal{M}'U(a)^*$? If this is nontrivial, one obtains a chiral component of a CFT (Guido-Longo-Wiesbrock '98). Integrable perturbation of CFT?



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Summary and open problems

- Some integrable QFT, including the sinh-Gordon model and other models with diagonal S-matrices (with CDD facctors), have been constructed in a mathematically satisfactory way.
- Some of them can be realized on the same Hilbert space as the free field, by twisting the observables in wedges, thus by "perturbing" the Heisenberg algebra (CFT).
- Complete the proof of modular nuclearity for nondiagonal S-matrices.
- Self-adjointness of the bound state operators (⇒ models with bound states).
- Study the lightlike intersection (\implies CFT from integrable models?).
- Wedge-algebras by free product (Longo-T.-Ueda '17), local observables?
- Nets on the de Sitter spacetime (cf. Barata-Jäkel-Mund), conserved charges, proof of integrability?