The Principle of Least Action in Hamiltonian dynamics, Analysis and Symplectic geometry

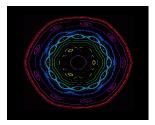
Alfonso Sorrentino

12th September 2014

Study of the dynamics of Hamiltonian systems

Order (stability) versus Chaos (instability)

- Methods from classical mechanics
- Perturbative methods (KAM theory,...)

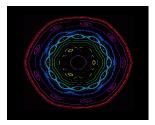


- Variational methods (Aubry-Mather theory)
- Geometric methods (Symplectic geometry, Floer Homology,...)
- PDE methods (Hamilton-Jacobi, weak KAM theory, ...)

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Aubry - Mather theory

Variational methods based on the *Principle of Least Lagrangian Action* (*"Nature is thrifty in all its actions"*, Pierre Louis Moreau de Maupertuis, 1744).

- Serge Aubry & John Mather '80s: twist maps of the annulus;
- John Mather '90s: Hamiltonian flows of *Tonelli* type.

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Tonelli Hamiltonian

Let M be finite dimensional closed Riemannian manifold. $H \in C^2(T^*M, \mathbb{R})$ is said to be *Tonelli* if:

- *H* is strictly convex in each fibre: $\partial_{pp}^2 H(x, p) > 0$;
- *H* is super-linear in each fibre:

$$\lim_{\|p\|\to+\infty}\frac{H(x,p)}{\|p\|}=+\infty \quad \text{uniformly in } x.$$

• Geodesic Flow

Let g be a Riemannian metric on M. The Hamiltonian (or Kinetic energy)

$$H(x,p) = \frac{1}{2} \|p\|_x^2 := \frac{1}{2} g_x(p,p)$$

corresponds to the geodesic flow on M.

• Hamiltonians from classical mechanics (Kinetic Energy + Potential Energy):

$$H(x,p) = \frac{1}{2} \|p\|_x^2 + U(x)$$

where $U: M \to \mathbb{R}$ represents the potential energy.

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$$H(x,p) = \frac{1}{2} ||p||_x^2 + U(x)$$

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The dynamics associated to any vector field X on M can be embedded into the flow of a Tonelli Hamiltonian $H: T^*M \longrightarrow \mathbb{R}$:

$$H(x,p) = \frac{1}{2} ||p||_x^2 + p \cdot X(x).$$

Lagrangian formalism

Let $H : T^*M \longrightarrow \mathbb{R}$ a Tonelli Hamiltonian. We can associate to it the so-called Lagrangian function $L : TM \longrightarrow \mathbb{R}$, where

$$L(x,v) := \sup_{p \in T_x^*M} (p \cdot v - H(x,p))$$

Euler-Lagrange equations: $\frac{d}{dt}\frac{\partial L}{\partial v} = \frac{\partial L}{\partial x} \longrightarrow$ Euler-Lagrange flow.

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Euler-Lagrange equations: $\frac{d}{dt}\frac{\partial L}{\partial v} = \frac{\partial L}{\partial x} \longrightarrow$ Euler-Lagrange flow. The Hamiltonian flow and the Euler-Lagrange flow are equivalent from a dynamical system point of view:

$$TM \xrightarrow{\Phi_t^L} TM$$

$$\mathcal{L} \qquad \qquad \downarrow \mathcal{L}$$

$$T^*M \xrightarrow{\Phi_t^H} T^*M$$

The Euler-Lagrange flow has an interesting variational characterization in terms of the Lagrangian Action Functional. If $\gamma : [a, b] \longrightarrow M$ is an abs. cont. curve, we define its action as:

$$A_L(\gamma) := \int_a^b L(\gamma(t), \dot{\gamma}(t)) dt.$$

 γ is a solution of the Euler-Lagrange flow if and only if it is an extremal for the fixed-end variational problem.

These extremals are not necessarily minimisers, although they are *local* minimisers, *i.e.* for very short times.

Example: In the geodesic flow case not all geodesics are length minimising! But (global) minimising geodesics do exist.

Questions:

- Do global minimisers exist?
- What are their dynamical/geometric properties?
- Does this minimising property translate into some rigid structure?

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Aubry-Mather theory

Idea:

Study orbits and invariant probability measures that minimise the Lagrangian action of \boldsymbol{L}

Aubry-Mather theory

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Study orbits and invariant probability measures that minimise the Lagrangian action of $L - \eta_c(x) \cdot v$, where η_c is any smooth closed 1-form on M with cohomology class c.

Observation:

- η_c closed $\implies L$ and $L \eta_c$ have the same Euler-Lagrange flow.
- The corresponding Hamiltonian is $H(x, \eta_c(x) + p)$.

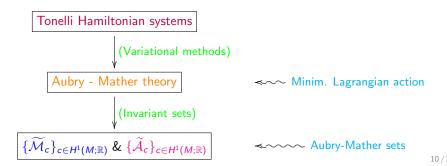
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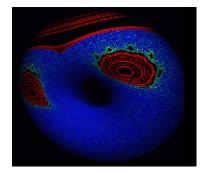
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Aubry-Mather sets

The Aubry-Mather sets are:

- non-empty and compact;
- invariant under the Hamiltonian flow;
- *supported* on Lipschitz graphs (Mather's graph theorem);

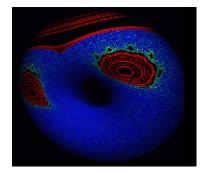


(Credits to Dr. Oliver Knill, Harvard)

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In particular:

- Symplectic invariant \longleftrightarrow Invariant under the action of simplectomorphisms.
- Lagrangian Structure \longleftrightarrow The are supported on Lipschitz Lagrangian, graphs of cohomology class c.

Proposition

If Λ is an invariant Lagrangian graph in $(T^*M, \omega_{\text{stand.}})$ with cohomology class c, then all orbits on Λ (resp. invariant prob. measures supported on Λ), minimises the action of $L - \eta_c(x) \cdot v$, where η_c is any smooth closed 1-form on M, with cohomology class c.

Therefore:

Aubry-Mather sets \leftrightarrow Invariant Lagrangian graphs (when they exist)

Invariant Lagrangian graphs are very rare. What if they do not exist?

Link to PDEs & Symplectic geometry

Proposition

If Λ is an invariant Lagrangian graph in $(T^*M, \omega_{\text{stand.}})$ with cohomology class c, then $\Lambda = \text{Graph}(c + du)$ and

$$H(x, c + du(x)) = \alpha(c).$$

Therefore, one can study viscosity solutions and subsolutions of Hamilton-Jacobi equations:

Aubry-Mather sets ↔ supported on the "graphs" of the differentials of these weak solutions (Uniqueness set)

Weak KAM theoryHomogenization of Hamilton-Jacobi equation(Albert Fathi '90s)(à la Lions-Papanicolaou-Varadhan and Evans)

For a fixed *c*, the sets M_c and A_c lie in an energy level $\{H(x, p) = \alpha(c)\}$. The function

$$\alpha: H^1(M; \mathbb{R}) \longrightarrow \mathbb{R}$$

is what we call Minimal average action or Effective Hamiltonian.

• It corresponds to the *Homogenized Hamiltonian*.

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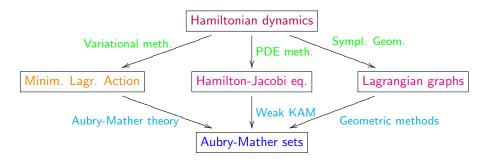
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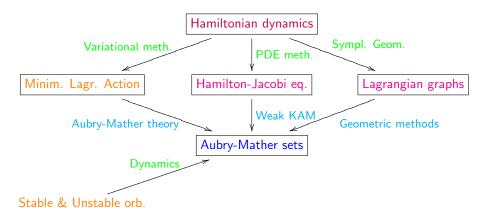
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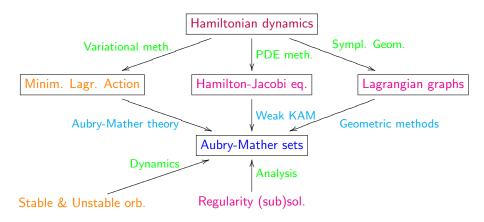
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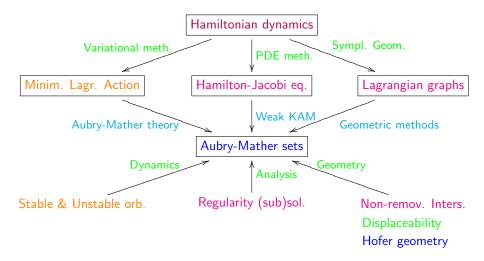
This function is also related to what is called Symplectic shape:

- α(c) = inf{k : the sublevel {H(x, p) ≤ k} contains Lagrangian graphs of cohomology class c}.
- α is related to Hofer geometry on the group of Hamiltonian diffeomorphisms (Sorrentino-Viterbo, Geom&Top 2010)









Research Interests

Structure of these Action-minimizing sets and H-J equation:

- Generic topological properties. Symplectic and contact properties.
- Implications to dynamics and Symplectic geometry;
- Implications to the regularity of viscosity solutions and subsolutions of Hamilton-Jacobi equation;
- Generalised forms of Homogenization of Hamilton-Jacobi equation.

Properties of the minimal average action:

- Symplectic properties and relation to Hofer geometry.
- Regularity, lack of regularity and geometric/dynamical implications.

Birkhoff Billiards (proposal for a SIR project 2014):

- Rigidity phenomena; (Length) Spectral properties.
- Integrability and Birkhoff conjecture.

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Some results on I

- A. S., *On the total disconnectedness of the quotient Aubry set*, Ergodic Theory Dynam. Systems 28 (2008), Vol. 1.
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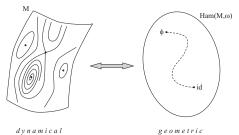
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- A. S. and Claude Viterbo Action minimizing properties and distances on the group of Hamiltonian diffeomorphisms, Geom. & Topol. 14 (2010), no. 4.
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Relation to Hofer geometry

Consider the group of (compactly supported) Hamiltonian diffeomorphism $Ham(M, \omega)$:

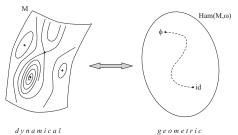


Fundamental (unsolved) question

What is the relation between the geometry of this curve and the dynamics of the system?

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Minimal average action \longleftrightarrow Asymptotic distance from Identity

Research Projects

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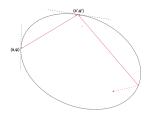
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Playing with Billiards

What is a (mathematical) billiard?



The billiard ball moves on a rectilinear path: when it hits the boundary it reflects elastically according to the standard reflection law:

angle of reflection = angle of incidence.

This is a conceptually simple model, yet mathematically very complicated,

In collaborations with Vadim Kaloshin (University of Maryland, USA) we have an ongoing project aimed at studying two important (and related) questions:

- Is it possible to hear the shape of a billiard? Can a planar convex domain be characterized in terms of the lengths of its periodic orbits, i.e., its Length spectrum (or Marked length spectrum), as conjectured by Guillemin and Melrose?
- Birkhoff conjecture on the integrability of convex billiards. Namely: the only integrable billiards are billiards in circles and ellipses.



Thank you for your attention!