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Spin-orbit coupling: the effect of rheology

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- This work is part of a postdoctoral project from the São Paulo Research Foundation (FAPESP - Grant 2021/11306-0)
 - ▶ Title: Spin-orbit coupling: the effect of the rheology
 - ▶ Supervisor: Clodoaldo Ragazzo (IME/USP)
 - ▶ Collaboration: Alexandre Correia (U. Coimbra)
- Our main objective is to study the scenario for which spin-orbit coupling occurs, eventually considering the effect of the body's rheology on this phenomenon
- This research is in its early stages. We will discuss a few preliminary results regarding the rotational dynamics of Hyperion, a moon of Saturn that is currently rotating chaotically, without considering the body's rheology
- Open-source software being developed at <https://github.com/vitor-de-oliveira/spin-orbit/tree/dev>

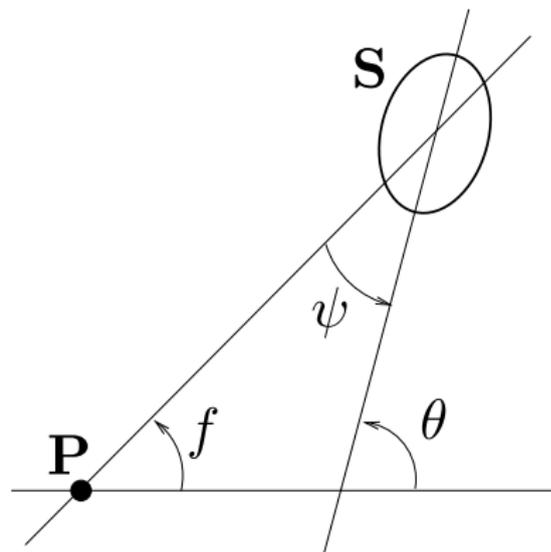
Physical model

S: Satellite

P: Planet

f : true anomaly

θ : rotation angle



Assumptions

- S orbits P in a fixed Keplerian ellipse with semi-major axis a , eccentricity e , and instantaneous radius r
- spin axis parallel to the largest principal moment of inertia and perpendicular to the orbit plane
- the only forces that act on S are the ones generated by the gravitational field of P

Spin-orbit resonance (SOR)

If T is the body's orbital period, then a spin-orbit resonance of type p/q , for two relatively prime integers p, q , with $q > 0$, is a solution for the system such that

$$\theta(t + Tq) = \theta(t) + 2\pi p \quad \forall t \in \mathbb{R},$$

i.e., after q revolutions the orbiting body has made p rotations about its spin axis

SOR in the Solar System

- Earth's Moon
 - ▶ 1/1 resonance
- Pluto and Charon
 - ▶ 1/1 resonance (both)
- Mercury
 - ▶ 3/2 resonance
- Many other moons (Phobos, Deimos, Io, Europa, Ganymede, Callisto, ...)
 - ▶ 1/1 resonance

- $I_1 < I_2 < I_3$: principal moments of inertia of S
- I_3 : moment of inertia about the spin axis

Rotational dynamics of an almost rigid body

$$I_3 \ddot{\theta} = -\frac{3}{2}(I_2 - I_1) \frac{Gm_P}{r^3} \sin 2(\theta - f) - 3k_2 \frac{Gm_P^2 R^5}{a^6} \tau (\bar{L}\dot{\theta} - \bar{N})$$

where

$$\bar{L}(e) = \frac{1}{(1 - e^2)^{9/2}} \left(1 + 3e^2 + \frac{3}{8}e^4 \right)$$

$$\bar{N}(e) = \frac{1}{(1 - e^2)^6} \left(1 + \frac{15}{2}e^2 + \frac{45}{8}e^4 + \frac{5}{16}e^6 \right)$$

Physically relevant parameters

orbital eccentricity

e

equatorial flattening

$$\gamma := \frac{I_2 - I_1}{I_3}$$

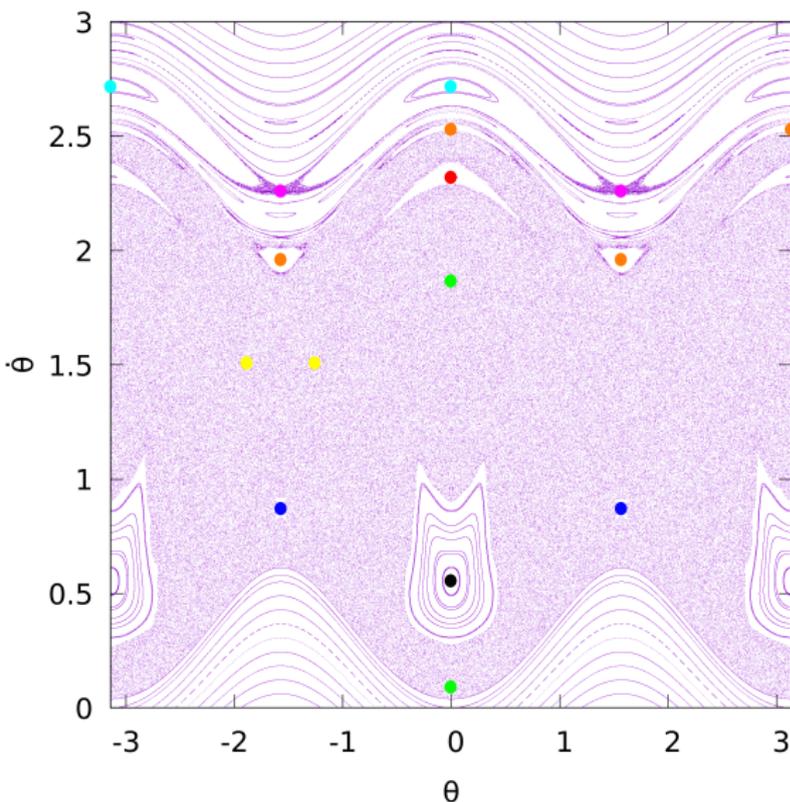
dissipation constant

$$K := 3k_2 \frac{Gm_P^2 R^5}{a^6} \tau$$

Hyperion

- Moon of Saturn
- Chaotic rotation
- Very aspherical shape, being nearly twice as long as it is across (Voyager 2)
- Physical parameters: $e \approx 0.1$ and $\gamma \approx 0.264$
- Main references: Wisdom and Peale (1984) and Wisdom (1987)

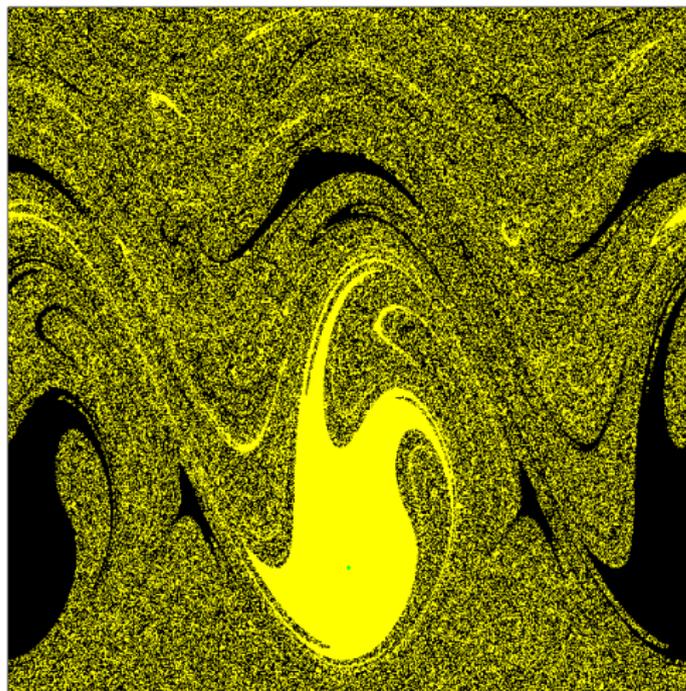
Phase space – $\theta \times \dot{\theta}$



Resonances:

- 1/1 stable (black)
- 2/1 stable (red)
- 1/2 stable (blue)
- 2/2 unstable (green)
- 3/2 unstable (yellow)
- 5/2 stable (light blue)
- 5/2 unstable (pink)
- 9/4 stable (orange)

Basin of attraction for resonance 1/1 with $K = 10^{-2}$

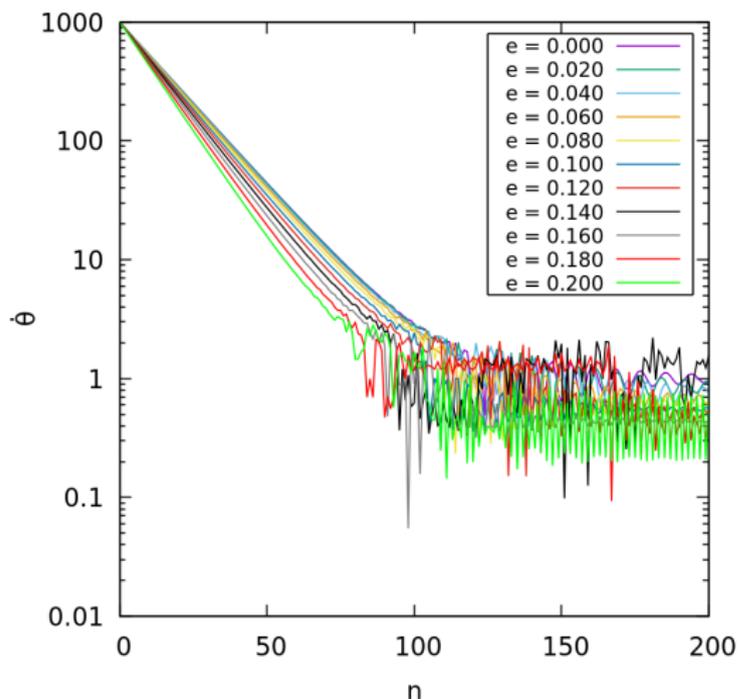


The basin of attraction of the stable 1/1 SOR presents a very complex structure

Hyperion - Phase space varying e (gif)

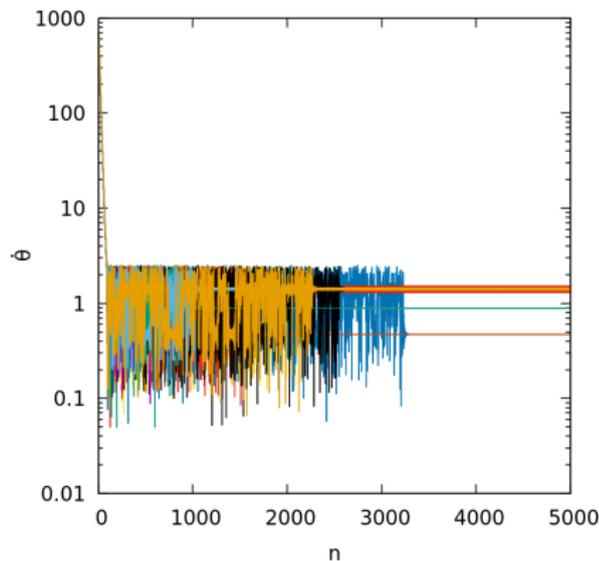
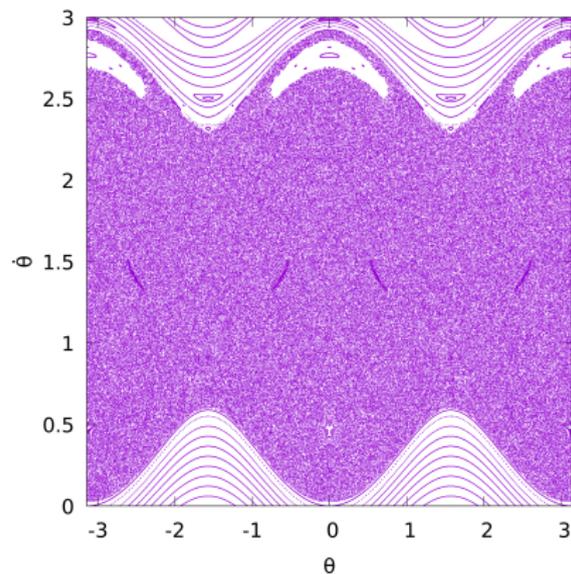
Hyperion - Phase space varying γ (gif)

Time series of $\dot{\theta}(n)$



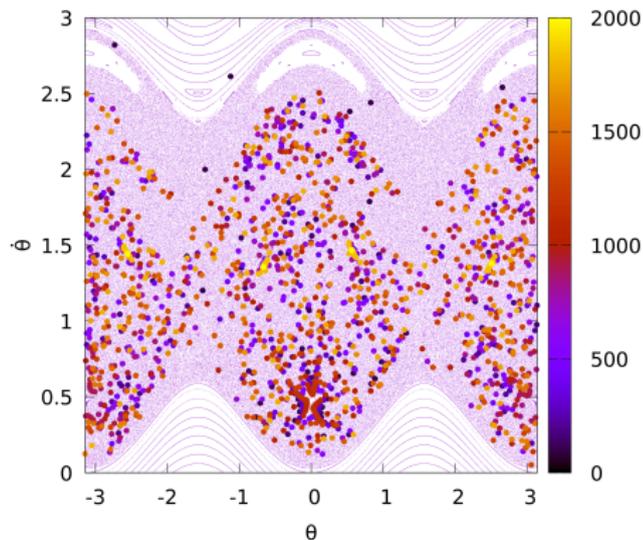
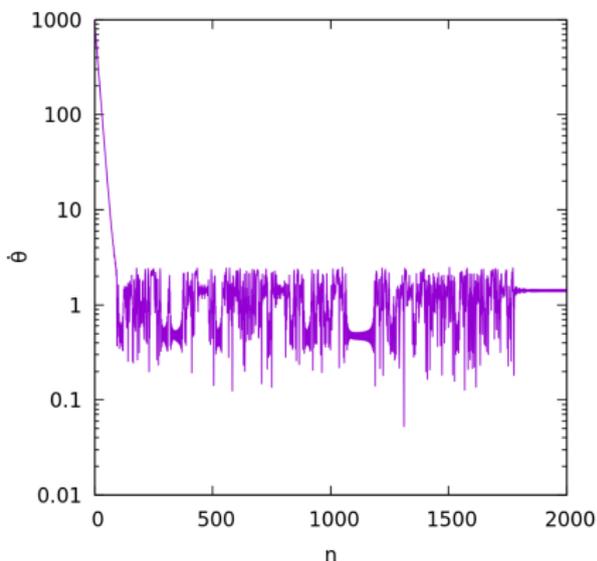
Same initial condition for $K = 10^{-2}$ and different values of the eccentricity e . In all cases, the angular velocity decays exponentially at first, with slope depending on e , before experimenting the chaotic region and finally converging to a SOR, which also depends on e .

Hyperion for $e = 0.14$



(left panel) The phase space for $e = 0.14$ presents a large chaotic region. (right panel) 20 initial conditions that are very close to each other converge to 4 different final states.

Metastable state



(left panel) Time series of $\dot{\theta}$ for a solution with $e = 0.14$. (right panel) Same orbit on phase space. The colorbox represent the orbit cycle n . We observe that the trajectory spends some time around the 1/1 SOR before converging to a higher-order resonance.

Next steps

- Measure the size of the basin of attraction for different spin-orbit resonances and different values of γ and e
- Account for changes in the equatorial flattening by considering a linear viscoelastic rheology model for the body, and analyze how these changes affect the entrapment into spin-orbit resonances and the timescale of metastable states

Main references

- Celletti, A. and Chierchia, L., *CM&DA*, 2008
- Celletti, A., *ZAMP*, 1990
- Goldreich P. and Peale S., *The Astronomical Journal*, 1966
- Ragazzo, C. and Ruiz, L. S., *CM&DA*, 2017