

Birkhoff Orbital Elements in Restricted 3- and 4-Body Problems

CELMEC VIII

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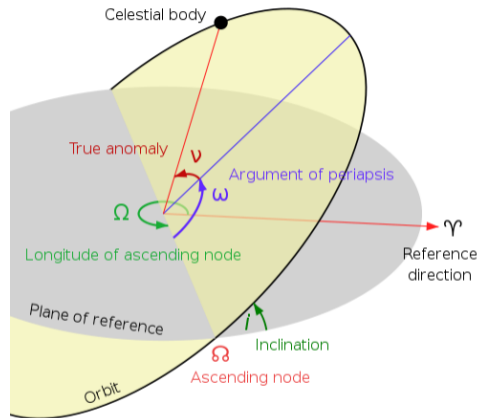
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Motivation

- Relative 2-Body Problem is an **integrable** dynamical system
- Emits the Keplerian orbital elements via integrals of motion
- Perturbations (J_2 , SRP, etc.) yield time-varying or *osculating* elements
- In N -body problems, $N \geq 3$, the Keplerian orbital elements are **no longer well-defined**



Keplerian Orbital Elements

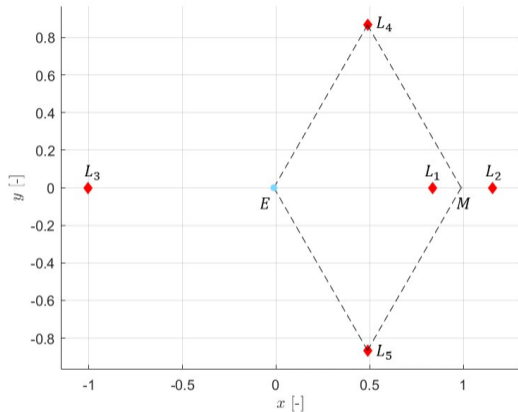


Restricted 3-Body Problem

- Restricted 3-Body Problem (R3BP) is an autonomous, **nearly integrable** Hamiltonian dynamical system

$$H = \frac{1}{2}(p_x^2 + p_y^2 + p_z^2) + yp_x - xp_y - \frac{1-\mu}{r_1} - \frac{\mu}{r_2}$$

- R3BP emits special solutions:
 - Equilibria
 - Periodic orbits
 - Quasi-periodic orbits
 - Hyperbolic invariant manifolds



R3BP synodical rotating frame



Birkhoff orbital elements are a **local analogue** to orbital elements in nearly integrable Hamiltonian dynamical systems, e.g. R3BP, defined using **action-angle coordinates** in a **Birkhoff normal form** about a bounded invariant manifold.

In order to compute Birkhoff orbital elements, we must first:

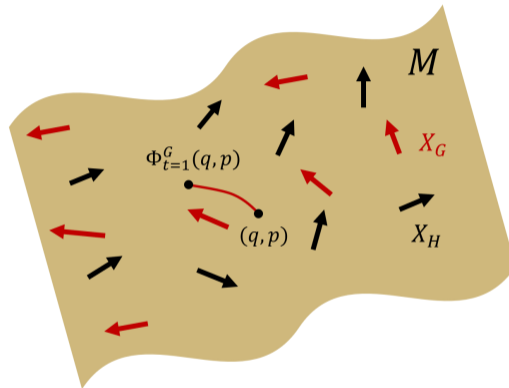
- Re-center about bounded invariant manifold
 - Eq. point, periodic orbit, reducible torus
- Diagonalize quadratic terms H_2 , e.g. $H_2 = \lambda q_1 p_1 + \frac{\omega_p}{2}(q_2^2 + p_2^2) + \frac{\omega_v}{2}(q_3^2 + p_3^2)$
- Complexify H_2 , e.g. $H_2 = \lambda q_1 p_1 + i\omega_p q_2 p_2 + i\omega_v q_3 p_3$
- Apply Lie series method (Jorba 1997, Rosales 2021)



Lie series method

- Maintain integrability to higher order via Lie series
- Uses Lie transforms, i.e., the time one flow of X_G , defined by a generating function G
- Obtain Birkhoff orbital elements via **action-angle** transformations in a **Birkhoff normal form**

$$\hat{H} = \sum_{k=2}^N \hat{H}_k(I) + R_N(q, p)$$

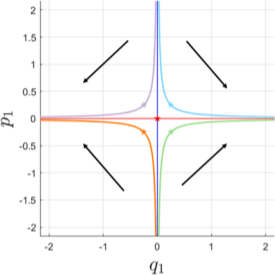
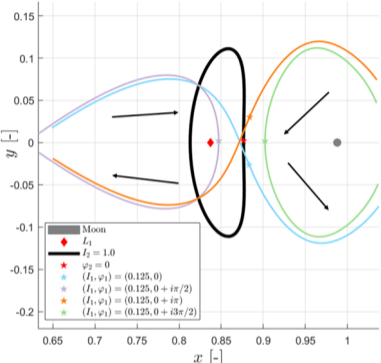
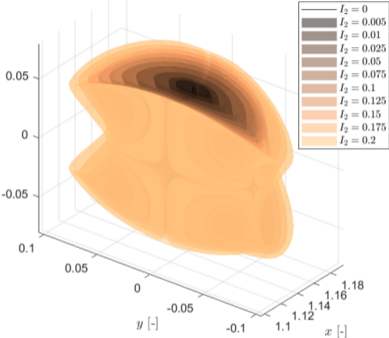


Birkhoff OEs: $(x, y, z, \dot{x}, \dot{y}, \dot{z}) \mapsto (I_s, \varphi_s, I_p, \varphi_p, I_v, \varphi_v)$



Examples in EM R3BP

L_2 Lissajous Orbits - Variable I_2 , Fixed $I_3 = 0.1$



Left: Lissajous orbits at a fixed vertical action, I_3

Right: Saddle trajectory about a periodic orbit with representation in saddle plane.

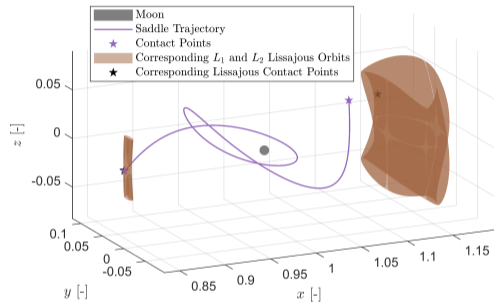


An Existence Theorem

Theorem

For each point (q, p) in the phase space of the Restricted 3-Body Problem, there exists a time $t^ \in \mathbb{R}$ such that $\phi_{t^*}(q, p)$ has a Birkhoff orbital element representation.*

We say that **each point has a Birkhoff orbital element representation** via the Birkhoff normal form flow backwards time $-t^*$. However, this representation is **non-unique**, as showed in the figure...

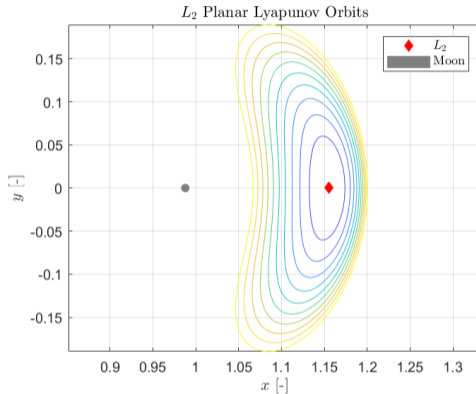


Example of non-unique representation in Birkhoff orbital elements in R3BP



Periodic Orbits in R3BP

- Libration point periodic orbits have Birkhoff orbital element representations
- We can also re-center our coordinates about an orbit outside this region
 - Halo orbits, DROs, etc.
- We can apply this same ideology to periodic orbits in R4BPs
 - HR4BP, BCP, and QBCP
 - Dynamically equivalents of equilibria
 - Resonant orbits



*L*₂ Planar Lyapunov Orbits via Birkhoff
Orbital Elements



Guaranteeing Symplectic Re-centering

Re-centering about periodic orbits in the R3BP and R4BP requires:

Theorem (Symplectic Floquet Theorem)

Let

$$H(\theta, X, Y) = X^T A(\theta) X + X^T B(\theta) Y + Y^T C(\theta) Y$$

be an n -d.o.f. Hamiltonian, where $A(\theta), B(\theta), C(\theta) \in \mathcal{C}^0(\mathbb{T}^1, M_n\mathbb{C})$, and $A(\theta), C(\theta)$ are symmetric for all θ . Then, \exists a symplectic, linear and periodic time-dependent change of variables that transforms H into a new autonomous Hamiltonian H_F on variables (x, y) :

$$H_F(x, y) = x^T \hat{A} x + x^T \hat{B} y + y^T \hat{C} y$$

and with $\hat{A}, \hat{B}, \hat{C} \in M_n\mathbb{C}$, and \hat{A}, \hat{C} symmetric matrices.



Summary

- **Birkhoff orbital elements** are a **local analogue** to orbital elements in the R3BP and HR4BP (resp. BCP, QBCP) defined using **action-angle coordinates** in a **Birkhoff normal form** centered about a bounded invariant manifold
- Every point in phase space can be represented by *at least* one set of Birkhoff orbital elements in the R3BP and R4BP

Future Work

- Finish building computational tool for periodic perturbation
- Study reducibility of tori to apply to quasi-periodic invariant tori

