

Noncommutative geometry of CP_q^ℓ

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The algebra of antiholomorphic forms on the quantum projective space \mathbb{CP}_q^ℓ is constructed. The associated Dolbeault-Dirac operator defines a 0-dimensional $U_q(\mathfrak{su}(\ell + 1))$ -equivariant spectral triple. (Based on [F. D'Andrea, L. Dąbrowski, arXiv:0901.4735]).

'QG' match 'NCG'

QG is q-deformed matrix group à la Woronowicz (and dual QUEA), or its homogeneous space.

NCG is à la Connes:

riemannian, spin manifold \longleftrightarrow spectral triple $(\mathcal{A}, \mathcal{H}, D)$,

where \mathcal{A} is a $*$ -algebra, \mathcal{H} is a Hilbert space representation of \mathcal{A} and D is an operator on \mathcal{H} s.t. $D = D^\dagger$, $(D - z)^{-1} \in \mathcal{K}$ and $[D, \mathcal{A}] \subset B(\mathcal{H})$.

(Here compact M & unital \mathcal{A} only).

Even if $\exists \gamma$ s.t. $\gamma^2 = 1$, $a\gamma = \gamma a$, $\forall a \in \mathcal{A}$, $D\gamma = -\gamma D$; so $\mathcal{H} = \mathcal{H}_+ + \mathcal{H}_-$.

Prototype: given a spin manifold M with a Riemannian metric g :

$\mathcal{A} = C^\infty(M, \mathbb{C})$ (commutative), $\mathcal{H} = L^2(\Sigma, \text{vol}_g)$, and $D = \not{D}$.

Even ($\gamma = \gamma_5$) iff $\dim M$ is even.

It satisfies further 7 conditions which permit reconstruction [Connes08].

Two crucial properties directly pass to $\text{NC}(\text{diff})G$ of Connes

$(D - z)^{-1} \in \mathcal{K} \Rightarrow D|_{\mathcal{H}_\pm}$ coupled to gauge connections, vector bundles, projections, or their classes in the K-theory, are Fredholm
 \rightsquigarrow 'topological' invariants via index computations.

Instead $[D, \mathcal{A}] \subset B(\mathcal{H})$ serves for 'locality'.

The tension between them + (some of) 7 other 'axioms', Connes formulated for NC algebras, permit spectacular calculations.

Among the axioms is *dimension* given by the asymptotics of $\text{spec} D$.

Another one is *reality*, c.f. the talk by G. Dossena.

- How to construct NC spectral triples (S.T.) on QG ?

Our strategy: equivariance under QUEA and 'harmonic analysis'.

E.g. 'isospectral' D , c.f. [CP03a, DLS05, DLPS05, DDLW07].

But on the standard Podleś quantum sphere $S_{q0}^2 \equiv \mathbb{CP}_q^1$, \exists also D [DS03] with exponential spectrum (and a 0^+ -summable S.T.)

This D has elegant description [SW04] in terms of $U_q(\mathfrak{su}(2))$ which act as derivations on S_{q0}^2 , and so $[D, a]$ are bounded.

Along this line, [K04] constructed D on q-flag manifolds, e.g. \mathbb{CP}_q^ℓ .

However, $\text{spec} D$ & $(D-z)^{-1} \in \mathcal{K}$ are not addressed.

We completed this task [DD09]: constructed 'Dirac-Dolbeault' operator D with derivation property (thus $[D, a] \subset B(\mathcal{H})$) and by relating D^2 to a Casimir element control $\text{spec} D$.

It is exponentially growing (so $(D-z)^{-1} \in \mathcal{K}$) and yields a 0^+ -dim ST.

I'll recall \mathbb{CP}_q^1 , describe \mathbb{CP}_q^2 (only spin^c) [DDL08-2] and mention \mathbb{CP}_q^ℓ .

- $\mathbb{CP}_q^1 = SU_q(2)/U(1) = S_q^3/U(1), \quad 0 < q < 1.$

The symmetry Hopf $*$ -algebra is $U_q(\mathfrak{su}(2))$, generated by K, K^{-1}, E, F

$$KE = qEK, \quad [E, F] = (q - q^{-1})^{-1}(K^2 - K^{-2})$$

(+ coproduct, counit and antipode).

The dual $\mathcal{A}(SU_q(2))$ is a $U_q(\mathfrak{su}(2))$ -bimodule $*$ -algebra w.r.t. \triangleright and \triangleleft .

Let $\forall N \in \mathbb{Z}$,

$$\Gamma_N = \left\{ a \in \mathcal{A}(SU_q(2)) \mid a \triangleleft K = q^{\frac{N}{2}} a \right\}.$$

$\mathcal{A} := \Gamma_0$ is just the $*$ -algebra of $\mathbb{CP}_q^1 = S_{q0}^2$.

Each Γ_N is a \mathcal{A} -bimodule. Also, a $U_q(\mathfrak{su}(2))_{\triangleright}$ -module and as such

$$\Gamma_N \simeq \bigoplus_{n-|N| \in 2\mathbb{N}} V_n, \quad (1)$$

where V_n is the $\text{spin}_{\frac{1}{2}n}$ irrep & \simeq is unitary wrt the Haar state on Γ_N .

The Casimir

$$\mathcal{C}_q = (q - q^{-1})^{-2} (q^{\frac{1}{2}}K - q^{-\frac{1}{2}}K^{-1})^2 + FE$$

has spectrum

$$\mathcal{C}_q|_{V_n} = [\frac{n+1}{2}]^2 \cdot id$$

with multiplicity $\dim V_n = n + 1$. (Here $[x] := \frac{q^x - q^{-x}}{q - q^{-1}}$).

Antiholomorphic forms: $\Omega = \Omega^0 \oplus \Omega^1$, w/ $\Omega^0 = \mathcal{A}$ and $\Omega^1 = \Gamma_2$.

The Dolbeault operator and its Hermitian conjugate are derivations

$$\begin{aligned} \bar{\partial} : \Omega^0 &\rightarrow \Omega^1, & a &\mapsto a \lhd F, \\ \bar{\partial}^\dagger : \Omega^1 &\rightarrow \Omega^0, & a &\mapsto a \lhd E. \end{aligned}$$

The Dolbeault-Dirac operator D on $\Omega^0 \oplus \Omega^1$ is

$$D(\omega_0, \omega_1) := (\omega_1 \lrcorner E, \omega_0 \lrcorner F) ,$$

and satisfies

$$D^2\omega = \omega \lrcorner (\mathcal{C}_q - [\tfrac{1}{2}]^2).$$

Using $\lrcorner \mathcal{C}_q = \mathcal{C}_q \lrcorner$, the decomposition (1) of Γ_N , and

$$D\gamma + \gamma D = 0, \quad \text{where} \quad \gamma = 1 \oplus -1,$$

$$\text{spec} D = \{0, \pm \sqrt{[k][k+1]} \mid k \in \mathbb{N} + 1\}$$

with multiplicity 1 (constant 0-forms) and $2k + 1$, resp.

Real spinors? Tensor Ω^\bullet with Γ_{-1} ($=\sqrt{\text{canonical bundle}}$):

$$\mathcal{H} = (\Omega^0 \oplus \Omega^1) \otimes_{\mathcal{A}} \Gamma_{-1} \simeq \Gamma_{-1} \oplus \Gamma_1 .$$

Then the *real* \not{D} is D twisted with the Grassmann connection of Γ_{-1} ; it agrees [SW04] w/ [DS03] and \exists a real structure J ($=$ spin structure).

Byproduct: compute the Dolbeault cohomology $H_{\bar{\partial}}^{\bullet}(\mathbb{CP}_q^2)$.

Define *harmonic* forms $\mathfrak{H}^n = \{\omega \in \Omega^{(0,n)} \mid D\omega = 0 (\equiv \bar{\partial}\omega = \bar{\partial}^{\dagger}\omega = 0)\}$.

Prop.[hodge] For all n , there is an orthogonal decomposition

$$\Omega^{(0,n)} = \mathfrak{H}^n \oplus \bar{\partial}\Omega^{(0,n-1)} \oplus \bar{\partial}^{\dagger}\Omega^{(0,n+1)} . \quad (2)$$

$\exists!$ harmonic form in each cohomology class:

$$H_{\bar{\partial}}^n(\mathbb{CP}_q^1) \simeq \mathfrak{H}^n = \ker D \Big|_{\Omega^{(0,n)}} .$$

Pf. As classically.

Corr.

$$H_{\bar{\partial}}^0(\mathbb{CP}_q^1) = \mathbb{C} , \quad H_{\bar{\partial}}^1(\mathbb{CP}_q^1) = 0 .$$

- $\mathbb{CP}_q(2)$

The symmetry Hopf \ast -algebra $U_q(\mathfrak{su}(3))$, $0 < q < 1$.

Generators K_i, K_i^{-1}, E_i, F_i , $i = 1, 2$, with $K_i = K_i^*, F_i = E_i^*$,
& relations

$$\begin{aligned} [K_i, K_j] &= 0, \quad K_i E_i K_i^{-1} = q E_i, \quad [E_i, F_i] = (q - q^{-1})^{-1} (K_i^2 - K_i^{-2}) \\ K_i E_j K_i^{-1} &= q^{-1/2} E_j, \quad [E_i, F_j] = 0, \quad \text{if } i \neq j, \end{aligned}$$

+ (Serre)

$$E_i^2 E_j - (q + q^{-1}) E_i E_j E_i + E_j E_i^2 = 0 \quad \forall i \neq j. \quad (3)$$

With the q -commutator $[a, b]_q := ab - q^{-1}ba$, (3) can be rewritten as

$$[E_i, [E_j, E_i]_q]_q = 0, \quad \text{or,} \quad [[E_i, E_j]_q, E_i]_q = 0.$$

(+ coproduct, counit and antipode).

*-irreps of $U_q(\mathfrak{su}(3))$ are classified by $n_1, n_2 \in \mathbb{N}^2$, have dim.

$$\leadsto \frac{1}{2}(n_1 + 1)(n_2 + 1)(n_1 + n_2 + 2), \quad (4)$$

and act on $V(n_1, n_2)$ with on. basis $|n_1, n_2, j_1, j_2, m\rangle$, where

$$j_i = 0, 1, 2, \dots, n_i, \quad \frac{j_1 + j_2}{2} - |m| \in \mathbb{N}, \quad (5)$$

as (explicit) multiplication (K 's) or weighted shift operators (E, F 's).
(The h.w.v. with weight $(q^{n_1/2}, q^{n_2/2})$ being $|n_1, n_2, n_1, 0, \frac{1}{2}n_1\rangle$).

\exists Casimir operator [R-P91] (add $H := (K_1 K_2^{-1})^{2/3}$ and H^{-1})

$$\begin{aligned} \mathcal{C}_q = (q - q^{-1})^{-2} & \left((H + H^{-1}) \left\{ (qK_1 K_2)^2 + (qK_1 K_2)^{-2} \right\} + H^2 + H^{-2} - 6 \right) \\ & + (qH K_2^2 + q^{-1} H^{-1} K_2^{-2}) F_1 E_1 + (qH^{-1} K_1^2 + q^{-1} H K_1^{-2}) F_2 E_2 \\ & + qH [F_2, F_1]_q [E_1, E_2]_q + qH^{-1} [F_1, F_2]_q [E_2, E_1]_q \end{aligned} \quad (6)$$

$\mathcal{C}_q = \mathcal{C}_q^*$ is central and $\mathcal{C}_q|_{V(n_1, n_2)}$ is

$$\leadsto \left[\frac{1}{3}(n_1 - n_2) \right]^2 + \left[\frac{1}{3}(2n_1 + n_2) + 1 \right]^2 + \left[\frac{1}{3}(n_1 + 2n_2) + 1 \right]^2 \times \text{id}. \quad (7)$$

$\mathcal{A}(SU_q(3))$ is generated by 9 elements u_j^i ($i, j = 1, \dots, 3$) & relations

$$\begin{aligned} u_k^i u_k^j &= q u_k^j u_k^i, & u_i^k u_j^k &= q u_j^k u_i^k, & \forall i < j, \\ [u_l^i, u_k^j] &= 0, & [u_k^i, u_l^j] &= (q - q^{-1}) u_l^i u_k^j, & \forall i < j, k < l, \end{aligned}$$

$$\sum_{\pi \in S_3} (-q)^{|\pi|} u_{\pi(1)}^1 u_{\pi(2)}^2 u_{\pi(3)}^3 = 1.$$

(+ *-structure, coproduct, counit and antipode).

$\mathcal{A}(SU_q(3))$ is a $U_q(\mathfrak{su}(3))$ -bimodule *-algebra, which we turn into a left $U_q(\mathfrak{su}(3)) \triangleright \otimes U_q(\mathfrak{su}(3)) \blacktriangleright$ -module *-algebra. As such, by Peter-Weyl

$$\mathcal{A}(SU_q(3)) \simeq \bigoplus_{(n_1, n_2) \in \mathbb{N}^2} V_{(n_1, n_2)} \otimes V_{(n_1, n_2)},$$

where \simeq is isometry wrt. the Haar state.

Important:

$$\mathcal{C}_q \triangleright = \mathcal{C}_q \blacktriangleright. \tag{8}$$

View [Meyer], [Welk] $\mathbb{CP}_q^2 = S_q^5/U(1)$ and in turn $S_q^5 = SU_q(3)/SU_q(2)$.
Dually,

$$\mathcal{A}(S_q^5) := \left\{ a \in \mathcal{A}(SU_q(3)) \mid h \blacktriangleright a = \epsilon(h)a, \forall h \in U_q(\mathfrak{su}(2)) \right\}$$

(the $*$ -subalgebra generated by $z_i := u_i^3$ and [VS] relations), where $U_q(\mathfrak{su}(2))$ is generated by $\{K_1, K_1^{-1}, E_1, F_1\}$.

$\mathcal{A}(S_q^5)$ is a module over $(U_q(\mathfrak{su}(2)) \blacktriangleright)' = U_q(\mathfrak{su}(3)) \triangleright \otimes U_q(\mathfrak{u}(1)) \blacktriangleright$,
where $U_q(\mathfrak{u}(1))$ is generated by $\{K_1 K_2^2, K_1^{-1} K_2^{-2}\}$, and decomposes

$$\mathcal{A}(S_q^5) \simeq \bigoplus_{(n_1, n_2) \in \mathbb{N}^2} V_{(n_1, n_2)} ,$$

where

$$K_1 K_2^2 \blacktriangleright |_{V_{(n_1, n_2)}} = q^{n_2 - n_1} \text{ id.}$$

In turn

$$\begin{aligned}\mathcal{A}(\mathbb{CP}_q^2) &:= \left\{ a \in \mathcal{A}(S_q^5) \mid K_1 K_2^2 \blacktriangleright a = a \right\} \\ &= \left\{ a \in \mathcal{A}(SU_q(3)) \mid h \blacktriangleright a = \epsilon(h)a, \ \forall h \in U_q(\mathfrak{u}(2)) \right\},\end{aligned}$$

where $U_q(\mathfrak{u}(2))$ is generated by $U_q(\mathfrak{su}(2))$ and $U_q(\mathfrak{u}(1))$.

The $*$ -algebra $\mathcal{A}(\mathbb{CP}_q^2)$ is generated by $p_{ij} := (u_i^3)^* u_j^3$, $(p_{ij})^* = p_{ji}$,
 $+$ relations $P^2 = P$, $P = P^*$, $\text{tr}_q(P) := q^4 p_{11} + q^2 p_{22} + p_{33} = 1$,
in terms of $\{P\}_{ij} = p_{ij}$.

$\mathcal{A}(\mathbb{CP}_q^2)$ is an $U_q(\mathfrak{su}(3))\blacktriangleright$ -module and decomposes

$$\mathcal{A}(\mathbb{CP}_q^2) \simeq \bigoplus_{n \in \mathbb{N}} V_{(n,n)}.$$

The Dolbeault complex

Think of $\mathcal{A}(\mathbb{CP}_q^2) \hookrightarrow \mathcal{A}(S_q^5)$ as a quantum principal $U(1)$ -bundle, and sections of associated line bundle over \mathbb{CP}_q^2 as $U(1)$ -equivariant maps. For $N \in \mathbb{Z}$, define $\mathcal{A}(\mathbb{CP}_q^2)$ -bimodules

$$\mathcal{L}_N := \left\{ a \in \mathcal{A}(S_q^5) \mid K_1 K_2^2 \blacktriangleright a = q^N a \right\}; \quad (9)$$

or \equiv , being $\mathcal{A}(S_q^5) \hookrightarrow \mathcal{A}(SU_q(3))$ a quantum principal $SU_q(2)$ -bundle,

$$\mathcal{L}_N = \left\{ a \in \mathcal{A}(SU_q(3)) \mid K_1 K_2^2 \blacktriangleright a = q^N a, \ h \blacktriangleright a = \epsilon(h)a, \ \forall h \in U_q(\mathfrak{su}(2)) \right\}.$$

($\mathcal{L}_0 = \mathcal{A}(\mathbb{CP}_q^2)$).

Each \mathcal{L}_N is $\mathcal{A}(\mathbb{CP}_q^2)$ -bimodule and also a left $U_q(\mathfrak{su}(3))$ -module and decomposes

$$\mathcal{L}_N \simeq \bigoplus_{n \in \mathbb{N}} V_{(n, n+N)}.$$

Motivated by the case $q = 1$ [GS99] (Kähler structure of \mathbb{CP}^2) let

$$\Omega^{(0,0)} := \mathcal{L}_0 = \mathcal{A}(\mathbb{CP}_q^2) , \quad \Omega^{(0,2)} := \mathcal{L}_3 .$$

Instead using the spin $1/2$ *-irrep $\tau : U_q(\mathfrak{su}(2)) \rightarrow \text{Mat}_2(\mathbb{C})$

$$\tau(K_1) = \begin{pmatrix} q^{1/2} & 0 \\ 0 & q^{-1/2} \end{pmatrix} , \quad \tau(E_1) = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} ,$$

we define $\Omega^{(0,1)}$ as the $\mathcal{A}(\mathbb{CP}_q^2)$ -bimodule

$$\begin{aligned} \Omega^{(0,1)} := & \left\{ v = (v_+, v_-) \in \mathcal{A}(SU_q(3))^2 \mid \right. \\ & \left. \mid K_1 K_2^2 \blacktriangleright v = q^{\frac{3}{2}} v, (h_{(1)} \blacktriangleright v) \tau(S(h_{(2)})) = \epsilon(h) v, \forall h \in U_q(\mathfrak{su}(2)) \right\} . \quad (10) \end{aligned}$$

It is not difficult to see that $\Omega^{(0,1)}$ has $U_q(\mathfrak{su}(3))\blacktriangleright$ -decomposition

$$\Omega^{(0,1)} \simeq \bigoplus_{n \geq 1} V_{(n,n)} \oplus \bigoplus_{n \geq 0} V_{(n,n+3)} .$$

Now Dolbeault cochain and dual chain complexes.

Prop. Set $X := F_2 F_1 - 2[2]^{-1} F_1 F_2$, $Y := E_2 E_1 - 2[2]^{-1} E_1 E_2$.

The maps

$$\bar{\partial} : \Omega^{(0,0)} \rightarrow \Omega^{(0,1)}, \quad \bar{\partial} a := (X^* \blacktriangleright a, E_2 \blacktriangleright a), \quad (11a)$$

$$\bar{\partial} : \Omega^{(0,1)} \rightarrow \Omega^{(0,2)}, \quad \bar{\partial} v := -E_2 \blacktriangleright v_+ - Y \blacktriangleright v_-, \quad (11b)$$

with $v = (v_+, v_-)$, are well defined and $\bar{\partial}^2 = 0$. Similarly, the maps

$$\bar{\partial}^\dagger : \Omega^{(0,2)} \rightarrow \Omega^{(0,1)}, \quad \bar{\partial}^\dagger b := (-F_2 \blacktriangleright b, -Y^* \blacktriangleright b), \quad (11c)$$

$$\bar{\partial}^\dagger : \Omega^{(0,1)} \rightarrow \Omega^{(0,0)}, \quad \bar{\partial}^\dagger v := X \blacktriangleright v_+ + F_2 \blacktriangleright v_-, \quad (11d)$$

are well defined and $(\bar{\partial}^\dagger)^2 = 0$.

As in the commutative case, is $\Omega^{(0,\bullet)}$ a graded algebra?

Since $\Omega^{(0,1)}$ and $\Omega^{(0,2)}$ are $\Omega^{(0,0)}$ -bimodules and $\dim. \leq 2$, we know how to multiply, except two 1-forms $v = (v_+, v_-)$ and $w = (w_+, w_-)$, for which we define

$$v \wedge_q w := \frac{2}{[2]} (q^{\frac{1}{2}} v_+ w_- - q^{-\frac{1}{2}} v_- w_+).$$

It can be seen that $v \wedge_q w \in \Omega^{(0,2)} = \mathcal{L}_3$ and associativity is OK!

This makes $\Omega^{(0,\bullet)}$ a graded algebra.

Moreover, $\Omega^{(0,\bullet)}$ is a left $U_q(\mathfrak{su}(3))_\triangleright$ module $*$ -algebra.

Using the (faithful) Haar state φ on $SU_q(3)$ we define a non-degenerate inner product

$$\langle \omega_1, \omega_2 \rangle := \varphi(a_1^* a_2 + v_{1+}^* v_{2+} + v_{1-}^* v_{2-} + b_1^* b_2), \quad (12)$$

w.r.t. which $U_q(\mathfrak{su}(3))_\triangleright$ is unitary [DDL08] and the decomposition $\Omega^{(0,\bullet)} := \bigoplus_n \Omega^{(0,n)}$ is orthogonal.

The operators $\bar{\partial}$ and $\bar{\partial}^\dagger$, defined via \blacktriangleright commute with $U_q(\mathfrak{su}(3))_{\blacktriangleright}$. Moreover, from [DDL08] follows $\forall v$ with entries in $\mathcal{A}(SU_q(3))$ that $h^* \blacktriangleright v = (h \blacktriangleright)^\dagger v$, wrt. $\langle \cdot, \cdot \rangle$ and so $\bar{\partial}^\dagger = (\bar{\partial})^\dagger$.

Prop.[gd] $\bar{\partial}$ is graded-derivation: $\forall a, b \in \Omega^{(0,0)}, v \in \Omega^{(0,1)}, c \in \Omega^{(0,2)}$,
 $\bar{\partial}(ab) = a(\bar{\partial}b) + (\bar{\partial}a)b, \bar{\partial}(av) = (\bar{\partial}a) \wedge_q v + a(\bar{\partial}v), \bar{\partial}(va) = (\bar{\partial}v)a - v \wedge_q (\bar{\partial}a)$,
 while $\bar{\partial}^\dagger$ satisfies:

$$[\bar{\partial}^\dagger, a]v = 2[2]^{-1}(F_2 \blacktriangleright a)v_- + q(X \blacktriangleright a)v_+, [\bar{\partial}^\dagger, a]c = -q^{\frac{3}{2}}(F_2 \blacktriangleright a, F_1 F_2 \blacktriangleright a)c.$$

(Pf. omitted).

Hence $(\Omega^{(0,\bullet)}, \bar{\partial})$ gives a left-covariant differential calculus; of dim. 2 (since ‘antiholomorphic’).

The spectral triple

On a Kähler spin_c manifold \exists a ‘Dolbeault-Dirac’ operator $\bar{\partial} + \bar{\partial}^\dagger$.

On \mathbb{CP}_q^2 , we need wider generality, for $s \in \mathbb{R}^+$ (a parameter) let

$$D(a, v, b) := (\bar{\partial}^\dagger v, \bar{\partial}a + s\bar{\partial}^\dagger b, s\bar{\partial}v) . \quad (13)$$

Let $\mathcal{H}_+ := \overline{\Omega^{(0,0)} \oplus \Omega^{(0,2)}}^{\langle \cdot, \cdot \rangle}$ and $\mathcal{H}_- = \overline{\Omega^{(0,1)}}^{\langle \cdot, \cdot \rangle}$.

Let $\mathcal{H} := \mathcal{H}_+ \oplus \mathcal{H}_-$ with grading $\gamma := 1 \oplus -1$.

Let the spinor rep. of $\mathcal{A}(\mathbb{CP}_q^2)$ be just the left multiplication.

They have few good properties

Our spinor rep. is bounded since $\mathcal{A}(\mathbb{CP}_q^2) \subset \mathcal{A}(SU_q(3)) \subset \mathcal{A}(SU_q(3))^4$ and $\mathcal{H} \subset L^2(SU_q(3), \varphi)^4$, the left multiplication being bounded.

By Prop.[gd], $[D, a]$ acts via left multiplication by elements of $\mathcal{A}(SU_q(3))$,

$$[D, a]\omega = \left([\bar{\partial}^\dagger, a]v, [\bar{\partial}, a]b + s[\bar{\partial}^\dagger, a]c, s[\bar{\partial}, a]v \right), \quad \forall \omega = (b, v, c),$$

and is bounded for any $a \in \mathcal{A}(\mathbb{CP}_q^2)$.

Equivariance: since our $\mathcal{A}(\mathbb{CP}_q^2)$ -modules of forms are equivariant and the operators $\bar{\partial}$ and $\bar{\partial}^\dagger$ are $U_q(\mathfrak{su}(3))$ -invariant.

Still need: (essential) self-adjointness and $(D - z)^{-1} \in \mathcal{K}$?

We'll get both by diagonalizing D . For that

Lemma Fix $s = \sqrt{[2]}/2$, then

$$D^2\omega = [2]^{-1}(\mathcal{C}_q - 2) \blacktriangleright \omega .$$

Pf. Direct verification on each grade.

Lemma[spec] The $\ker D$ is the constant 0-forms, while non-zero eigenvalues, $\forall n \geq 1$, are

$$\begin{aligned} \pm \sqrt{\frac{2}{[2]}[n][n+2]} & \quad \text{with multiplicity } (n+1)^3 , \\ \pm \sqrt{[n+1][n+2]} & \quad \text{with multiplicity } \frac{1}{2}n(n+3)(2n+3) . \end{aligned}$$

Pf. Use $\mathcal{C}_q \blacktriangleright = \mathcal{C}_q \triangleright$ (8), decomposition of $\Omega^{(0,0)}$, $\Omega^{(0,1)}$, $\Omega^{(0,2)}$ into irreps of $U_q(\mathfrak{su}(3)) \triangleright +$ symmetry wrt. 0 ($D\gamma + \gamma D = 0$).

Since the eigenvalues of D grow exponentially, $(D + i)^{-\epsilon}$ is trace class $\forall \epsilon > 0$ and the metric dim is 0^+ , hence $(D - z)^{-1} \in \mathcal{K}$.

This justifies

Prop.[main] For $s = \sqrt{[2]/2}$ in (13), the datum $(\mathcal{A}(\mathbb{CP}_q^2), \mathcal{H}, D, \gamma)$ is a 0^+ -dimensional $U_q(\mathfrak{su}(3))$ -equivariant even spectral triple.

Byproduct:

As before $\mathfrak{H}^n = \ker D$ and Hodge decomposition works, and thus

$$H_{\bar{\partial}}^0(\mathbb{CP}_q^2) = \mathbb{C} , \quad H_{\bar{\partial}}^1(\mathbb{CP}_q^2) = H_{\bar{\partial}}^2(\mathbb{CP}_q^2) = 0 .$$

In [DD09] we have generalized all that to \mathbb{CP}_q^ℓ : with much more involved combinatorics of the Grassman algebra (cf. the talk by F. D'andrea).

The bounded comutators are again assured using special algebraic properties of D , and the compactness by relating D to an appropriate Casimir operator.

The main result is the construction of 0^+ -dimensional $U_q(\mathfrak{su}(\ell + 1))$ -equivariant even spectral triple.

If ℓ is odd and $N = \frac{1}{2}(\ell + 1)$, the spectral triple is a real with KO-dimension $2N \bmod 8$.

On the basis of the low dim cases and now the whole series \mathbb{CP}_q^ℓ , $\ell \in \mathbb{N}$, the works of Chakraborty-Pal and the fascinating paper [Neshveyev-Tuset07] on QG, built on the impressive collection of papers by Lusztig, Kazhdan and Etingof on analytic generalization of the formal equivalence by Drinfeld (via the twist and associators) between the the categories of $U_h(g)$ -modules and $U_g[[h]]$ -modules,

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I can declare NCG and QG as matched

Now starts the party:

should clear the isospectral versus exponential spectrum (or intermediate?)

establish such properties of QG representations (without their explicit form), that assure $[D, a] \in \mathcal{B}(\mathcal{H})$ and $|D|^{-1} \in K(\mathcal{H})$

work out the other conditions of Connes.

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