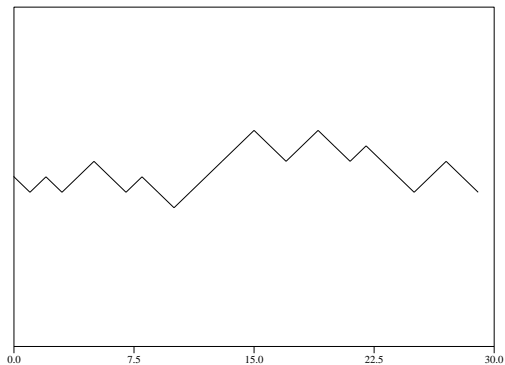


# RANDOM WALKS ON SOME NONCOMMUTATIVE SPACES

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Vietri Sul Mare, 01/09/2009

# Classical random walk



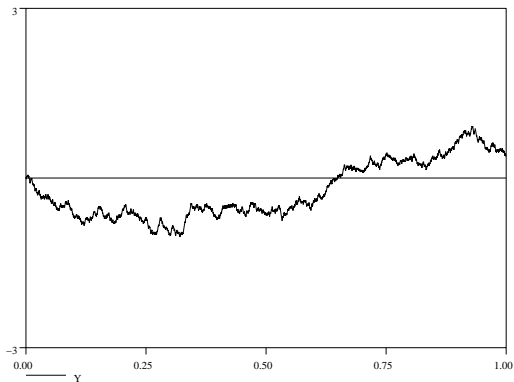
$$S_n = X_1 + \dots + X_n$$

$$X_k = \pm 1$$

# Brownian motion

Scale by  $\varepsilon$  in time and  $\sqrt{\varepsilon}$  in space.

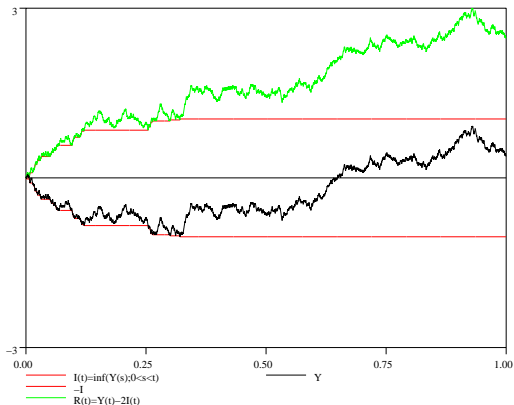
$$\sqrt{\varepsilon} \uparrow \rightarrow \varepsilon$$



# PITMAN THEOREM (1975)

$B_t; t \geq 0$  Brownian motion;  $I_t = \inf_{0 \leq s \leq t} B_s$

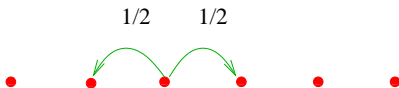
$R_t = B_t - 2I_t; t \geq 0$  is distributed as the norm of a three dimensional Brownian motion(=Bessel 3 process)



Explained by considering a random walk in a non-commutative space.

## DISCRETE VERSION

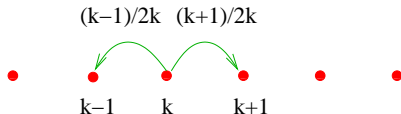
$$X_i = \pm 1; \quad S_n = X_1 + X_2 + \dots + X_n$$



$R_n = S_n - 2 \min_{0 \leq k \leq n} S_k$  is a Markov chain(=discrete Bessel 3 process)

$$P(R_{n+1} = k + 1 | R_n = k) = \frac{k + 1}{2k}$$

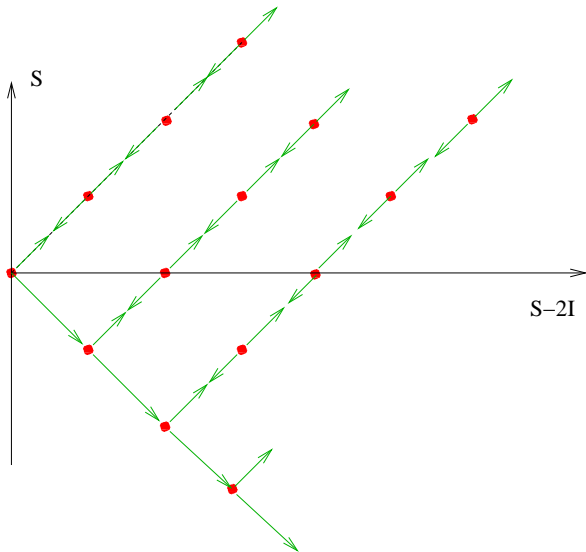
$$P(R_{n+1} = k - 1 | R_n = k) = \frac{k - 1}{2k}$$



when  $n \rightarrow \infty$   $S_{[nt]}/\sqrt{n} \rightarrow_{n \rightarrow \infty}$  Brownian motion

$R_{[nt]}/\sqrt{n} \rightarrow_{n \rightarrow \infty}$  norm of 3D-Brownian motion

# PROOF OF PITMAN'S THEOREM



## Quantum Bernoulli random walks

We "quantize" the set of increments of the random walk  $\{\pm 1\}$  to obtain  $M_2(\mathbf{C})$ .

The subset of hermitian operators in  $M_2(\mathbf{C})$  is a four dimensional real subspace, generated by the identity matrix  $I$  as well as the three matrices

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

The matrices  $\sigma_x, \sigma_y, \sigma_z$  are the Pauli matrices. They satisfy the commutation relations

$$[\sigma_x, \sigma_y] = 2i\sigma_z; \quad [\sigma_y, \sigma_z] = 2i\sigma_x; \quad [\sigma_z, \sigma_x] = 2i\sigma_y \quad (1)$$



## The random walk

For  $\omega$  a state on  $M_2(\mathbf{C})$ , in  $(M_2(\mathbb{C}), \omega)^{\otimes \mathbb{N}}$  we put

$$x_n = I^{\otimes(n-1)} \otimes \sigma_x \otimes I^{\otimes \infty}, \quad y_n = I^{\otimes(n-1)} \otimes \sigma_y \otimes I^{\otimes \infty}, \quad z_n = I^{\otimes(n-1)} \otimes \sigma_z \otimes I^{\otimes \infty}$$

$x_n$  is a commuting family of operators, a sequence of independent Bernoulli random variables.

$$X_n = \sum_{i=1}^n x_i; \quad Y_n = \sum_{i=1}^n y_i \quad Z_n = \sum_{i=1}^n z_i$$

are Bernoulli random walks.

They do not commute but obey

$$[X_n, Y_m] = 2iZ_{n \wedge m} \tag{2}$$

as well as the similar relations obtained by cyclic permutation of  $X, Y, Z$ .

$(X_n, Y_n, Z_n); n \geq 1$  is a *quantum Bernoulli random walk*.

## The spin process

Let  $S_n = \sqrt{I + X_n^2 + Y_n^2 + Z_n^2}$

**Proposition** For all  $n, m$  one has

$$[S_n, S_m] = 0$$

Thus we have a commutative process and we can try to compute its distribution.

## Theorem

Let  $\omega$  be the tracial state  $\frac{1}{2} Tr$ , then  $S_n$  is distributed as a Markov chain on the positive integers, with probability transitions

$$p(k, k+1) = \frac{k+1}{2k}; \quad p(k, k-1) = \frac{k-1}{2k}.$$

## Random walks on groups

$W$  = abelian group

$\hat{W}$  = dual group

$\xi \in \hat{W}$  = character of  $W$

$F(W)$  = algebra of functions on  $W$

$A(\hat{W})$  = group algebra of  $\hat{W}$

$$F(W) \rightarrow F(W \times W)$$

$$\Delta : A(\hat{W}) \rightarrow A(\hat{W}) \otimes A(\hat{W})$$

$$f(x) \rightarrow f(x + y)$$

$$\Delta(\xi) = \xi \otimes \xi$$

$$\mu : F(W) \rightarrow \mathbb{C}$$

$\phi$  = positive definite function on  $\hat{W}$

= probability measure on  $W$

$$\phi(\xi) = \int_W \xi(x) d\mu(x)$$

state  $\omega$  on  $A(\hat{W})$

$$\Omega = (W, \mu)^\infty$$

$$M = \otimes^\infty (A(\hat{W}), \omega)$$

$$Y_n = w_1 + \dots + w_n$$

$$f \rightarrow f(w_1 + \dots + w_n)$$

$$j_n : A(\hat{W}) \rightarrow M$$

$$j_{n+1} = (\Delta \otimes I^{\otimes(n+1)}) \circ I \otimes j_n$$

Markov operator

$$\Phi(f)(x) = \int_W f(x+y) d\mu(y) \quad \Phi(f) = (I \otimes \omega) \circ \Delta$$

## Random walks on duals of compact groups

Replace  $\hat{W}$  by a compact group  $G$ .

$\phi$  = continuous positive definite functions on  $G$ , with  $\phi(e) = 1$ .

= state  $\nu$  on  $A(G)$ .

$\nu$  = distribution of the increments.

$$\Phi_\nu : \mathcal{A}(G) \rightarrow \mathcal{A}(G)$$

$$\Phi_\nu = (I \otimes \nu) \circ \Delta$$

is a completely positive map. It generates a semigroup  $\Phi_\nu^n$ ;  $n \geq 1$ .

$$(\mathcal{N}, \omega) = (\mathcal{A}(G), \nu)^\infty$$

$j_n : \mathcal{A}(G) \rightarrow \mathcal{N}$  defined by  $j_n(\lambda_g) = \lambda_g^{\otimes n} \otimes I$

The morphisms  $(j_n)_{n \geq 0}$ , define a random walk on the noncommutative space dual to  $G$ , with Markov operator.

$$\Phi_\nu : \mathcal{A}(G) \rightarrow \mathcal{A}(G)$$

$$\Phi_\nu = (I \otimes \nu) \circ \Delta$$

The quantum Bernoulli random walk is obtained for  $G = SU(2)$ , and  $\nu$  the tracial state associated with the 2-dimensional representation.

## The dual of $SU(2)$ as a noncommutative space

$G = SU(2)$  = unitary  $2 \times 2$  matrices with determinant 1.

$$\hat{G} = \{1, 2, 3, \dots\}$$

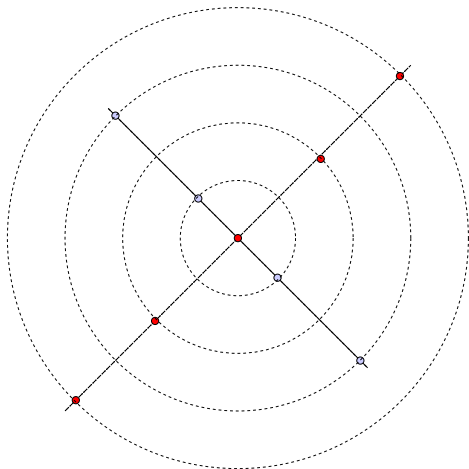
$$\mathcal{A}(SU(2)) = \oplus_{n=1}^{\infty} M_n(\mathbb{C})$$

is the noncommutative space dual to  $SU(2)$ .

The Pauli matrices belong to the Lie algebra  $\mathfrak{su}(2)$ , they define unbounded operators  $X, Y, Z$ , on  $L^2(SU(2))$ .

They generate oneparameter subgroups isomorphic to  $U(1)$ . This is true also of any linear combination  $xX + yY + zZ$  with  $x^2 + y^2 + z^2 = 1$ .





Noncommutative space underlying  $\mathcal{A}(SU(2))$

If you are in this space and measure your coordinate in some direction  $(x, y, z)$  using the operator  $xX + yY + zZ$ , and you will always find an integer.

You cannot measure coordinates in two different directions at the same time.

The operator  $D = \sqrt{I + X^2 + Y^2 + Z^2} - I$  is in the center of the algebra  $\hat{\mathcal{A}}(SU(2))$ , and therefore can be measured simultaneously with any other operator.

Its eigenvalues are the nonnegative integers  $0, 1, 2, \dots$ , and its spectral projections are the identity elements of the algebras  $M_n(\mathbb{C})$

$$D = \sum_{n=1}^{\infty} (n-1) I_{M_n(\mathbb{C})}$$

$M_n(\mathbb{C})$  is a kind of "noncommutative sphere of radius  $n-1$ ". Looking at the eigenvalues of the operators  $xX + yY + zZ$  the coordinate on this "radius" can only take the  $n+1$  values  $n, n-2, n-4, \dots, -n$ .

## Construction of the quantum Bernoulli random walk

$\omega$ =state on  $M_2(\mathbb{C})$ ,  $\nu = \omega^{\otimes \infty}$  on  $\mathcal{N} = \otimes_1^\infty M_2(\mathbb{C})$ .

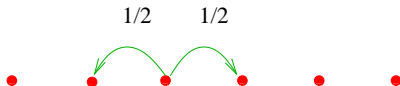
Construct morphisms  $j_n : \mathcal{A}(SU(2)) \rightarrow \mathcal{N}$  by

$$j_n(\lambda_g) = \rho_2(g)^{\otimes n} \otimes I^{\otimes \infty}$$

The family of morphisms  $(j_n)_{n \geq 1}$  is a stochastic noncommutative process with values in the dual of  $SU(2)$ .

This is just the iterated tensor product of the basic representation viewed as a random walk.

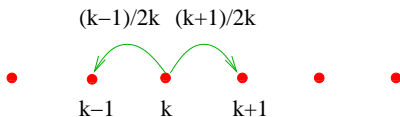
Restriction to a one parameter subgroup gives a Bernoulli random walk.

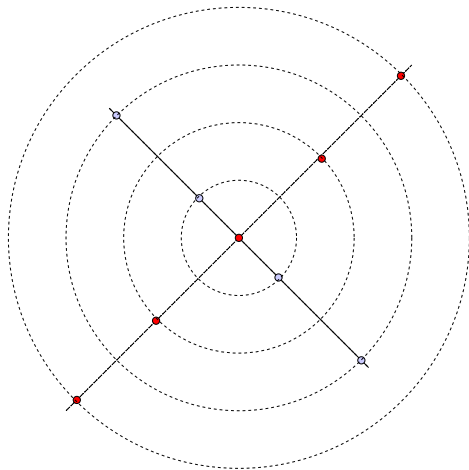


The spin process (radial part) is obtained by restriction of  $j_n$  to the center of the group algebra.

The restriction of the completely positive map  $\Phi$  to this center can be computed by the Clebsch Gordan formula

$$\rho_2 \otimes \rho_k = \rho_{k-1} \oplus \rho_{k+1}$$



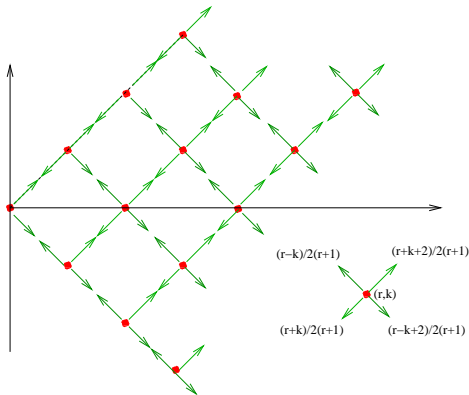


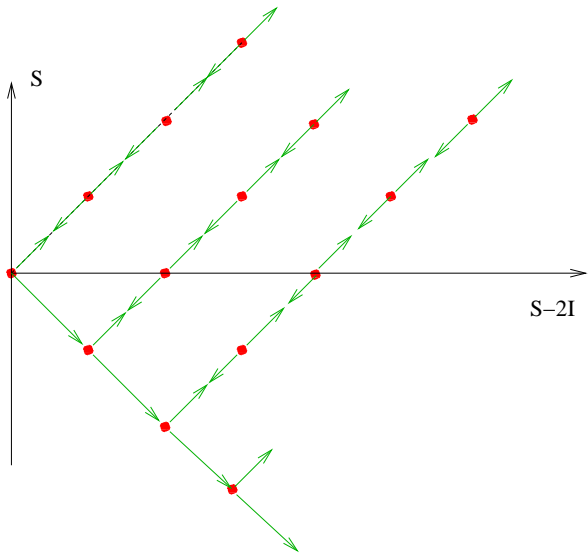
## RESTRICTION TO A MAXIMAL ABELIAN ALGEBRA

Restrict the Markov operator to the maximal abelian subalgebra generated by the center and a one parameter subgroup.

In the decomposition  $\mathcal{A}(SU(2)) = \oplus M_n(\mathbb{C})$  this is the algebra of diagonal operators.

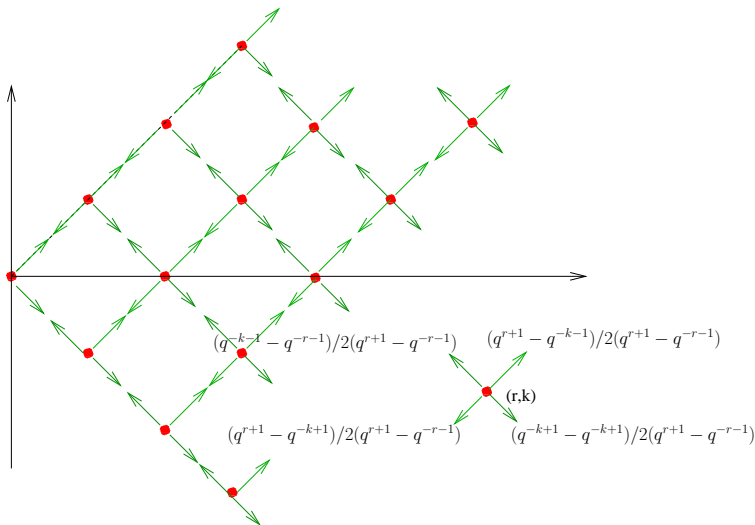
One gets probability transitions





# Kashiwara's crystallization

Replace  $SU(2)$  by Drinfeld/Jimbo/Woronowicz  $SU_q(2)$  then



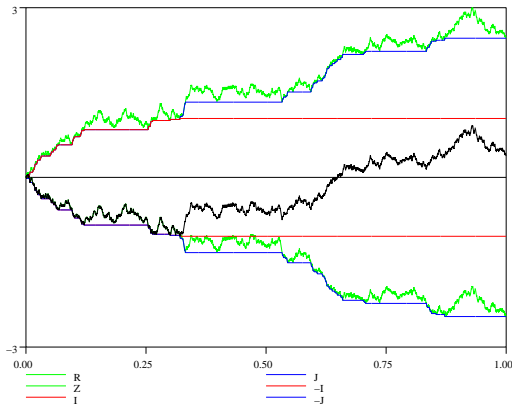
Let  $q \rightarrow 0$  then one obtains Pitman's theorem.



We can generalize the preceding construction to the quantum groups  $SU_q(n)$ .

# PITMAN OPERATORS

$$Y : [0, T] \rightarrow \mathbf{R}, \quad Y(0) = 0$$



$$PY(t) = Y(t) - 2 \inf_{0 \leq s \leq t} Y(s)$$

For all  $t$  one has  $PY(t) \geq 0$ , in particular  $PPY = PY$ .

## MULTIDIMENSIONAL PITMAN OPERATORS

$V$ =real vector space,  $\alpha \in V$ ,  $\alpha^\vee \in V^*$   $\alpha^\vee(\alpha) = 2$ .

$$P_\alpha Y(t) = Y(t) - \inf_{0 \leq s \leq t} \alpha^\vee(Y(s))\alpha$$

$$P_\alpha P_\alpha Y = P_\alpha Y$$

## Braid relations

If the angle between  $\alpha$  and  $\beta$  is  $\theta = \pi/n$  then

$$P_\alpha P_\beta P_\alpha \dots = P_\beta P_\alpha P_\beta \dots \quad (n \text{ terms})$$

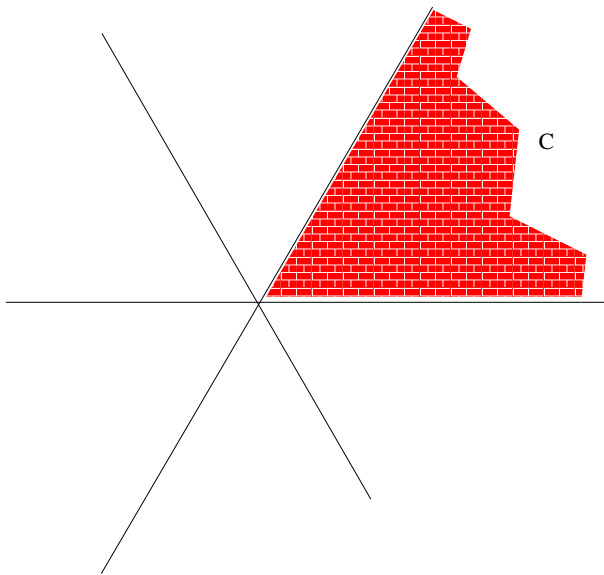
**Corollary:** Let  $(W, S)$  = Coxeter system on  $V$  and  $\alpha, \alpha^\vee$  = simple roots and coroots,

$C$  = Weyl chamber. To each  $s_\alpha \in S$  associate  $P_{s_\alpha}$ . For each  $w \in W$  with reduced decomposition  $w = s_{\alpha_1} \dots s_{\alpha_k}$  there exists

$$P_w = P_{s_{\alpha_1}} \dots P_{s_{\alpha_k}}$$

If  $w_0$  = longest element then  $P_{w_0} X$  takes values in  $C$ .

Example  $W = S_3$



## GENERALIZED PITMAN THEOREM

Let  $X$  be Brownian motion in  $V$   
then  $P_{w_0}X$  is Brownian motion "conditioned to stay in  $C$ ".

# DOOB'S CONDITIONED BROWNIAN MOTION

$$\psi(x) = \prod_{\beta \in R_+} \beta(x)$$

is a positive harmonic function on  $C$

$$p_t^W(x, y) = \sum_{w \in W} \varepsilon(w) p_t(x, w(y))$$

is the fundamental solution of Laplacian on  $W$  with Dirichlet boundary conditions

(=transition probabilities for Brownian motion killed at the boundary of  $C$ ).

$$q_t(x, y) = \frac{\psi(y)}{\psi(x)} p_t^W(x, y)$$

are the transition probabilities of Brownian motion conditioned to stay in  $C$ .

Fact:

when  $W = S_n$  (i.e. Weyl group of type  $A_{n-1}$ ) then Brownian motion conditioned to stay in  $C$  is the same as the motion of eigenvalues

$$(\lambda_1(t), \lambda_2(t), \dots, \lambda_n(t))$$

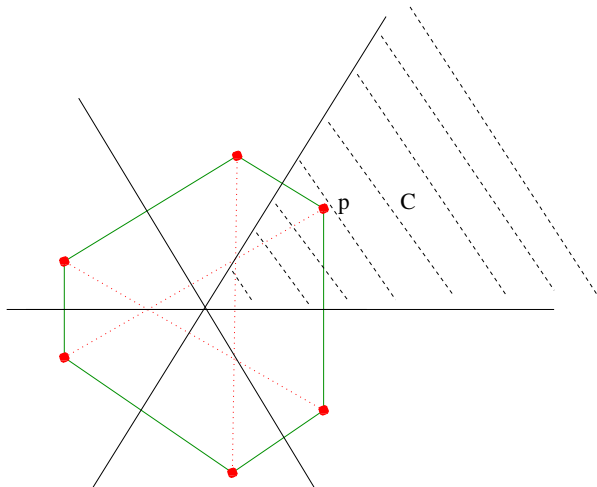
of a Brownian traceless hermitian matrix.

$$(M_{ij}(t))$$



# CONVERSE THEOREM

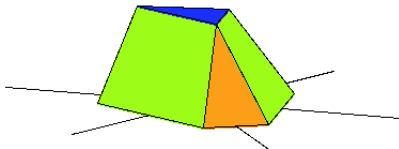
The conditional distribution of  $X(t)$  knowing  $P_{w_0}X(t) = p$  is the Duistermaat-Heckmann measure on the convex polytope with vertices  $w(p)$ ;  $w \in W$ .



Its Fourier transform is

$$\frac{1}{\prod_{\beta \in R} \beta(y)} \sum_{w \in W} \varepsilon(w) e^{i\langle p, y \rangle}$$

density is piecewise polynomial

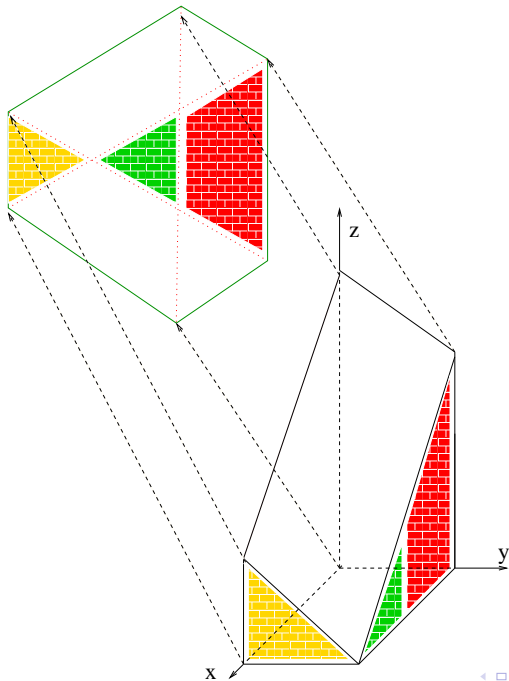


In order to recover  $X$  from  $P_{w_0}X$  we need a positive real number  $x_i$  for each  $s_i$  in  $P_{w_0} = P_{s_1} \dots P_{s_q}$ .

**Lemma** Given  $P_{w_0}X(t)$  the numbers  $(x_1, \dots, x_q)$  belong to a certain convex polytope. Their distribution is the normalized Lebesgue measure on this polytope.

Cristallographic case: Berenstein-Zelevinsky polytopes

The Duistermaat-Heckman measure is the image of this measure by an affine map.



$$0 < x < a$$

$$0 < y < b$$

$$0 < z < (a-x) + (b-y)$$