Overview	Classical Potential Theory	Dirichlet forms on von Neumann algebras	KMS-symmetry	Quantum Spins	Compact Quantum Groups	Lévy Process	HAP
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Noncommutative Potential Theory 1

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Themes.

- Review of Classical Potential Theory CPT
- Dirichlet forms on Standard Forms of von Neumann algebras
- KMS symmetric semigroups on C*-algebras
- Approach to equilibria in Quantum Spin Systems
- Quantum Lévy Processes on Compact Quantum Groups
- Characterization of Haagerup Approximation Property by Dirichlet forms

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Classical Potential Theory concerns properties of the Dirichlet integral

$$\mathcal{D}: L^2(\mathbb{R}^d, m) \to [0, +\infty] \qquad \mathcal{D}[u]:= \int_{\mathbb{R}^d} |\nabla u|^2 \, dm:$$

- lower semicontinuous quadratic form on the Hilbert space $L^2(\mathbb{R}^d, m)$
- finite on the Sobolev space $H^1(\mathbb{R}^d)$
- closed form of the Laplace operator

$$\Delta = -\sum_{k=1}^{d} \partial_k^2 \qquad \mathcal{D}[u] = \|\sqrt{\Delta}u\|_2^2$$

- generator of the heat semigroup $e^{-t\Delta}: L^2(\mathbb{R}^d, m) \to L^2(\mathbb{R}^d, m)$
- whose heat kernel

$$e^{-t\Delta}(x,y) = (4\pi t)^{-d/2} e^{-\frac{|x-y|^2}{4t}}$$

is the fundamental solution of the heat equation $\partial_t u + \Delta u = 0$

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The contraction property or Markovianity $\mathcal{D}[u \wedge 1] \leq \mathcal{D}[u]$ is responsible for

- Maximum Principle for solution of the Laplace equation $\Delta u = 0$
- Maximum Principle for solutions of the heat equation $\partial_t u + \Delta u = 0$
- contractivity, positivity preserving and continuity properties of the heat semigroup e^{-tΔ} on the spaces L²(ℝ^d, m), L[∞](ℝ^d, m), L¹(ℝ^d, m).

The Brownian motion (Ω, P_x, X_t) is the stochastic processes on \mathbb{R}^d associated to \mathcal{D}

$$(e^{-t\Delta}u)(x) = P_x(u \circ X_t)$$

whose polar sets *B* (avoided by the processes) are the Cap(B) = 0 sets for the electrostatic capacity associated to \mathcal{D} .

• The above properties are proved by the knowledge of the Green function

$$\Delta^{-1}u(x) = \int_{\mathbb{R}^d} G(x, y)u(y) m(dy) \qquad G(x, y) = |x - y|^{2-d} \qquad d \ge 3.$$

- Beurling and Deny (late '50) developed a kernel free potential theory generalizing the notion of Dirichlet integral to locally compact spaces.
- Fukushima (middle '60) achieved the construction of the associated Hunt process.

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Let $(\mathcal{M}, L^2(\mathcal{M}), L^2_+(\mathcal{M}), J)$ be a standard form of a von Neumann algebra \mathcal{M} . Let $\xi_0 \in L^2_+(\mathcal{M})$ be a fixed cyclic and separating vector and $\xi \wedge \xi_0 \in L^2_+(\mathcal{M})$ be the projection of a real vector $\xi = J\xi \in L^2(\mathcal{M})$ onto the positive cone $L^2_+(\mathcal{M})$.

Definition. (Dirichlet form)

A Dirichlet form $\mathcal{E}: L^2(\mathcal{M}) \to (-\infty, +\infty]$ is a l.s.c., quadratic form such that

- the domain $\mathcal{F} := \{\xi \in L^2(\mathcal{M}) : \mathcal{E}[\xi] < +\infty\}$ is dense in $L^2(\mathcal{M})$
- $\mathcal{E}[J\xi] = \mathcal{E}[\xi]$ real
- $\mathcal{E}[\xi \wedge \xi_0] \leq \mathcal{E}[\xi]$ Markovian
- $(\mathcal{E}, \mathcal{F})$ is a complete Dirichlet form if its matrix expansions for $n \ge 1$

$$\mathcal{E}_n[(\xi_{ij})_{ij}] := \sum_{ij} \mathcal{E}[\xi_{ij}]$$

are Dirichlet forms on $\mathcal{M} \otimes \mathbb{M}_n(\mathbb{C})$ (tacitly assumed since now on)

The domain \mathcal{F} is called *Dirichlet space* when endowed with the graph norm

$$\|\xi\|_{\mathcal{F}} := \sqrt{\mathcal{E}[\xi] + \|\xi\|_{L^2(\mathcal{M})}^2}$$

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Definition. (Markovian semigroup)

A self-adjoint C_0 -semigroup $\{T_t : t \ge 0\}$ on $L^2(\mathcal{M})$ is Markovian if

• $T_t J = J T_t$ $t \ge 0$

•
$$\xi \leq \xi_0 \quad \Rightarrow \quad T_t \xi \leq \xi_0 \quad t \geq 0$$

• $\{T_t : t \ge 0\}$ on $L^2(\mathcal{M})$ is completely Markovian if its matrix expansions

 $T_t^n([\xi_{ij}]_{ij}) := [T_t\xi_{ij}]_{ij}$

are Markovian semigroups on $L^2(\mathcal{M} \otimes \mathbb{M}_n(\mathbb{C}))$ (tacitly assumed since now on)

Consider the symmetric embedding $i_0 : \mathcal{M} \to L^2(\mathcal{M})$ $i_0(x) := \Delta_{\xi_0}^{1/4} x \xi_0$ and the faithful, normal state $\omega_0 : \mathcal{M} \to \mathbb{C}$ $\omega_0(x) := (\xi_0 | x \xi_0)_2$.

Theorem. (Modular ω_0 -symmetry)

Markovian semigroups are in 1:1 correspondence with C_0^* -continuous, positively preserving, contractive semigroups $\{S_t : t \ge 0\}$ on \mathcal{M} which are ω_0 -symmetric

$$\omega_0(S_t(x)\sigma_{-i/2}^{\omega_0}(y)) = \omega_0(\sigma_{-i/2}^{\omega_0}(x)S_t(y)) \qquad x, y \in \mathcal{M}_{\sigma^{\omega_0}}, \quad t > 0$$

through $i_0(S_t(x)) = T_t(i_0(x))$ $x \in \mathcal{M}$.

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Theorem. (Generalized Beurling-Deny correspondence)

Dirichlet forms are in 1:1 correspondence with Markovian semigroups by

$$\mathcal{E}[\xi] = \lim_{t \to 0} \frac{1}{t} (\xi | a - T_t \xi) \qquad a \in \mathcal{F}$$

or through the self-adjoint generator (L, dom(L))

$$T_t = e^{-tL} \qquad \mathcal{E}[a] = \left\|\sqrt{L}a\right\|_{L^2(A,\tau)}^2 \quad a \in \mathcal{F} = dom\left(\sqrt{L}\right).$$

In particular, Dirichlet forms are nonnegative $\mathcal{E} \ge 0$ and Markovian semigroups are positivity preserving and contractive.

• Extending Markovian semigroups from \mathcal{M} to $L^2(\mathcal{M})$ via non symmetric embeddings

$$i_{\alpha}(x) := \Delta^{\alpha}_{\xi_0} x \xi_0 \qquad \alpha \in [0, 1/2] \qquad \alpha \neq 1/4 \,,$$

produces semigroups on $L^2(\mathcal{M})$ which automatically commute with Δ_{ξ_0} .

 By duality and interpolation, Markovian semigroups extend to C₀-semigroups on noncommutative L^p(M) spaces, p ∈ [1, +∞).

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Theorem. (Ergodic Markovian semigroups)

The following properties are equivalent:

- the Markovian semigroup $\{T_t : t \ge 0\}$ on $L^2(\mathcal{M}, \omega)$ is ergodic: for $\xi, \eta \in L^2_+(\mathcal{M}, \omega)$ there exists t > 0 such that $(\xi | T_t \eta)_2 > 0$
- the Markovian semigroup {T_t : t ≥ 0} on L²(M, ω) is indecomposable: for some t > 0, T_t leaves invariant no proper face of the cone L²₊(M, ω)
- λ := inf{E[ξ] : ||ξ||₂ = 1} is a Perron-Frobenius eigenvalue:
 it is a simple eigenvalue with cyclic eigenvector ξ_λ ∈ L²₊(M,ω).

Faces *F* of the self-polar cone $L^2_+(\mathcal{M}, \omega)$ are in 1:1 correspondence with Peirce projections $P_e = eJeJ$ associated to projections $e \in \operatorname{Proj}(\mathcal{M})$

$$F = P_e(L^2_+(\mathcal{M},\omega)).$$

In the trace case, the above equivalences were established by L. Gross in his paper *Existence and uniqueness of physical ground states*, J. Funct. Anal. 10 (1972).

Let $\{\alpha_t : t \in \mathbb{R}\}$ be a strongly continuous automorphisms group on the C^{*}-algebra A, A_{α} the algebra of its analytic elements and let $\omega \in A_+^*$ be a KMS_{β}-state for $\beta \in \mathbb{R}$.

Definition. (KMS symmetric semigroups on C*-algebras)

A C_0 -semigroup $\{S_t : t \ge 0\}$ on A is KMS_β symmetric with respect to ω if

$$\omega(bS_t(a)) = \omega(\alpha_{-\frac{i\beta}{2}}(a)S_t(\alpha_{+\frac{i\beta}{2}}(b))) \qquad a, b \in B$$

for some dense, α -invariant, *-subalgebra $B \subseteq A_{\alpha}$.

• equivalently $\omega(\alpha_{-\frac{i\beta}{2}}(b)S_t(a)) = \omega(\alpha_{-\frac{i\beta}{2}}(a)S_t(b))$ $a, b \in B$

• KMS symmetry is a deformation of the KMS condition, in fact for t = 0 we get

$$\omega(ba)=\omega(\alpha_{-\frac{i\beta}{2}}(a)\alpha_{+\frac{i\beta}{2}}(b))=\omega(a\alpha_{+i\beta}(b)) \qquad a,b\in B$$

• In case $\{\alpha_t : t \in \mathbb{R}\}$ and $\{S_t : t \ge 0\}$ commute, KMS symmetry reduces to

$$\omega(bS_t(a)) = \omega(S_t(b)a) \qquad GNS \ symmetry$$

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also referred to as detailed balance.

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Proposition.

The following conditions are equivalent

- a C_0 -semigroup $\{S_t : t \ge 0\}$ on A is KMS_β symmetric with respect to ω
- for any a, b ∈ A and on the KMS-strip D_β ⊂ C there exists a bounded continuous function F_{a,b} : D_β → A, analytic in D_β such that for s ∈ R, t ≥ 0

 $F_{a,b}(s) = \omega(\alpha_{-s}(a)S_t(\alpha_{+s}(b))), \qquad F_{a,b}(s+i\beta) = \omega(\alpha_{+s}(b)S_t(\alpha_{-s}(a))).$

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Let $\omega_0 \in A^*_+$ be a KMS_{β}-state for $\{\alpha_t : t \in \mathbb{R}\} \subset Aut(A)$ and consider

- the cyclic GNS representation $(\pi_{\omega_0}, \mathcal{H}_{\omega_0}, \xi_0)$ of *A*
- the von Neumann algebra $\mathcal{M} := \pi_{\omega_0}(A)''$ acting on the space
- $L^2(\mathcal{M},\omega_0)\simeq \mathcal{H}_{\omega_0}$ carrying
- the standard form determined by $L^2_+(\mathcal{M},\omega_0) = \{\Delta^{1/4}_{\xi_0}\pi_{\omega_0}(A_+)\xi_0\}$
- the normal extension of ω_0 to \mathcal{M} given by $\omega_0(x) := (\xi_0 | \xi_0 x)_2$, $x \in \mathcal{M}$
- the modular automorphisms group $\{\sigma_t^{\omega_0} : t \in \mathbb{R}\}$ of \mathcal{M} .

Proposition.

- A KMS_{β} symmetric, C₀-semigroup {S_t : $t \ge 0$ } on A
 - leaves globally invariant the kernel of the cyclic representation: $S_t(\ker(\pi_{\omega_0})) \subseteq \ker(\pi_{\omega_0})$
 - extends to a ω₀-symmetric, C^{*}₀-semigroup {T_t : t ≥ 0} on the von Neumann algebra M by T_t π_{ω₀} = π_{ω₀} S_t
 - extends to a Markovian semigroup on $L^2(\mathcal{M}, \omega_0)$
 - determines a Dirichlet form on the standard form $(\mathcal{M}, L^2(\mathcal{M}, \omega_0), L^2_+(\mathcal{M}, \omega_0))$

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Example:	Bounded Dirichlet forms						

Example (C. JFA 147 (1997)).

On a standard form $(\mathcal{M}, L^2(\mathcal{M}), L^2_*(\mathcal{M}), J)$ consider j(x) := JxJ for $x \in \mathcal{M}$ and

- finite subsets $\{a_k : k = 1, \dots, n\} \subset \mathcal{M}, \{\mu_k, \nu_k : k = 1, \dots, n\} \subset (0, +\infty)$
- operators $d_k : L^2(\mathcal{M}) \to L^2(\mathcal{M})$ defined by $d_k := i(\mu_k a_k \nu_k j(a_k^*))$
- quadratic form on $L^2(\mathcal{M})$ given by $\mathcal{E}[\xi] := \sum_{k=1}^n \|d_k \xi\|_{L^2(\mathcal{M})}^2$

Then \mathcal{E} is

- *J*-real iff $\sum_{k=1}^{n} [\mu_k^2 a_k^* a_k \nu_k^2 a_k a_k^*] \in \mu \cap \mathcal{M}'$
- Markovian if moreover $\sum_{k=1}^{n} [\mu_{k}^{2} a_{k}^{*} a_{k} - \mu_{k} \nu_{k} (a_{k} j(a_{k}) + a_{k}^{*} j(a_{k}^{*})) + \nu_{k}^{2} a_{k} a_{k}^{*}] \xi_{0} \ge 0;$
- the associated Markovian semigroup is conservative, T_tξ₀ = ξ₀ for all t ≥ 0, if moreover the numbers (μ_k/ν_k)² are eigenvalues of the modular operator Δ_{ξ0} corresponding to eigenvectors a_kξ₀
- the generator has the form

$$L = \sum_{k=1}^{n} [\mu_k^2 a_k^* a_k - \mu_k \nu_k (a_k j(a_k) + a_k^* j(a_k^*)) + \nu_k^2 a_k a_k^*].$$

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Example (C.-Fagnola-Lindsay CMP 210 (2000)).

- Consider the canonical base $\{e_k : k \in \mathbb{N}\}$ of Hilbert space $h := l^2(\mathbb{N})$
- the C^{*}-algebra of compact operators $\mathcal{K}(h)$
- the von Neumann algebra of bounded operators $\mathcal{B}(h)$
- the Hilbert-Schmidt standard form $(\mathcal{B}(h), \mathcal{L}^2(h), \mathcal{L}^2_+(h), J)$
- fix parameters $\mu > \lambda > 0$ and set $\nu := (\lambda/\mu)^2$
- the state $\omega_{\nu}(x) := (1 \nu) \sum_{k \ge 0} \nu^{k} |e_{k}| > < e_{k}|$
- the cyclic vector $\xi_{\nu} := (1 \nu)^{1/2} \sum_{k \ge 0} \nu^{k/2} |e_k| > \langle e_k|$
- creation/annihilation operators $a^*(e_k) := \sqrt{k+1}e_{k+1}$ $a(e_k) := \sqrt{k}e_{k-1}$ $a(e_0) = 0$ satisfying the *Canonical Commutation Relation*: $aa^* - a^*a = I$.

Then the closure of the quadratic form $\mathcal{E}:\mathcal{L}^2(h) \to [0,+\infty)$

$$\mathcal{E}[\xi] := \|\mu a \xi - \lambda \xi a^* \|^2 + \|\mu a \xi^* - \lambda \xi^* a^* \|^2 \qquad \mathcal{F} := \text{linear span}\{|e_k > < e_l| : k, l \in \mathbb{N}\}$$

is a Dirichlet form and the associated Markovian semigroup reduces to an ergodic, Markovian, C₀-semigroup on $\mathcal{K}(h)$ leaving the state ω_{ν} invariant.

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- Consider the lattice \mathbb{Z}^d and the class \mathcal{L} of its finite subsets
- denote $A_X := \bigotimes_{x \in X} \mathbb{M}_2(\mathbb{C})$ the algebra of observables in $X \in \mathcal{L}$
- denote $A_0 = \bigcup_{X \in \mathcal{L}} A_X$ the normed algebra of local observables
- $A := \overline{A_0}$ the C*-algebra of quasi-local observabes
- consider an interaction $\Phi := \{ \Phi_X = \Phi_X^* \in A_X : X \in \mathcal{L} \}$

Then if $\lambda > 0$ is such that $\|\Phi\|_{\lambda} := \sup_{x \in \mathbb{Z}^d} \sum_{x \in X \in \mathcal{L}} |X| 4^{|X|} e^{\lambda \operatorname{diam}(X)} \|\Phi_X\| < +\infty$, a norm closable derivation *A* is defined on by

$$D(\delta):=A_0 \qquad \delta(a):=\sum_{X\cap Y
eq \emptyset} i[\Phi_Y,a] \qquad a\in A_X \quad X\in \mathcal{L}$$

and the automorphisms group $\{\alpha^{\Phi}_t : t \in \mathbb{R}\}$ generated by its closure satisfies

- analiticity: the evolution ℝ ∋ t → α^Φ_t(a) of local observables a ∈ A₀, extends analytically to the strip D_{βλ}, β_λ := ^λ/_{2||Φ||λ}
- finite group velocity: for $a \in A_{\{x\}}$, $b \in A_X$, $t \in \mathbb{R}$ we have

$$\|[\alpha_t^{\Phi}(a), b]\| \le 2\|a\| \cdot \|b\| \cdot |X| \cdot e^{-(\lambda \operatorname{dist}(x, X)) - 2|t|\Phi\|_{\lambda})}$$

Isotropic, anisotropic Heisenberg and Ising models correspond to different Φ .

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Example: Approach to equilibria of Quantum Spin Systems 1							

- Let $\omega \in A_+^*$ be a KMS_{β} state of the automorphisms group $\{\alpha_t^{\Phi} : t \in \mathbb{R}\}$
- consider the standard form (M, L²(M, ω), L²(M, ω)) of M := π_ω(A)" generated by the cyclic representation (π_ω, L²(M, ω), ξ_ω)
- consider the Pauli matrices $\{\sigma_j^x \in A_{\{x\}} : j = 0, 1, 2, 3\}$ at sites $x \in \mathbb{Z}^d$
- their images $a_j^x := \pi_\omega(\sigma_j^x) \in \mathcal{M}$
- denote $f_0 : \mathbb{R} \to \mathbb{R}$ the function $f_0(t) := (\cosh(2\pi t))^{-1}$

Theorem. (Y.M. Park, IDAQP Rel. Top. 3, (2000))

At sufficiently high temperature $\beta < \frac{\lambda}{\|\Phi\|_{\lambda}}$, the form $\mathcal{E} : L^2(\mathcal{M}, \omega) \to [0, +\infty]$

$$\mathcal{E}[\xi] := \sum_{x \in \mathbb{Z}^d} \sum_{j=0}^3 \mathcal{E}_{x,j}[\xi] \qquad \mathcal{E}_{x,j}[\xi] := \int_{\mathbb{R}} \|[\sigma_{t-i/4}(a_j^x) - j(\sigma_{t-i/4}(a_j^x))]\xi\|^2 f_0(t) dt$$

is a Dirichlet form with respect to the cyclic vector $\xi_{\omega} \in L^2_+(\mathcal{M}, \omega)$.

Proof: combines i) stability of the Markovian property and lower semicontinuity under superposition ii) the condition on the temperature implies that a_j^x are analytic elements for the modular group so that the forms $\mathcal{E}_{x,j}$ are well defined and also provide a dense domain in $L^2(\mathcal{M}, \omega)$ where \mathcal{E} is finite.

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Example:	Approach to equilibria of Qu	antum Spin Systems 1					

Theorem. (Y.M. Park J. Math. Physics 46 (2005))

The following properties are equivalent

- ω is an extremal KMS_{β} state for the automorphisms group { $\alpha_t^{\Phi} : t \in \mathbb{R}$ }
- ω is a factor state
- the Markovian semigroup $\{T_t : t \ge 0\}$ on $L^2(\mathcal{M}, \omega)$ is ergodic.

"Proof": by construction, $\{a_j^x : x \in \mathbb{Z}^d, j = 0, 1, 2, 3\}$ generates \mathcal{M} and one gets

$$\{\xi \in L^2(\mathcal{M},\omega): T_t\xi = \xi, t > 0\} = \overline{(\mathcal{M} \cap \mathcal{M}')\xi_0}.$$

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- Let $\omega \in A_+^*$ be a KMS_{β} state of the automorphisms group $\{\alpha_t^{\Phi} : t \in \mathbb{R}\}$
- consider the standard form (M, L²(M, ω), L²(M, ω)) of M := π_ω(A)["]
- consider the partial traces $\operatorname{Tr}_X : A \to A$ corresponding to $X \in \mathcal{L}$

Theorem. (A. Majewski-B. Zegarlinski Lett. Math. Phys. 36 (1996))

There exist $\lambda > 0$ such that $\|\Phi\|_{\lambda} < +\infty$ and $\beta > 0$ such that

• There exist $\gamma_X \in A_X$ normalized and rapidly decaying

$$\operatorname{Tr}_X(\gamma_X^*\gamma_X) = 1 \qquad \|\gamma_{X+j} - \operatorname{Tr}_i(\gamma_{X+j})\|_A \le c \cdot (1+|i-j|)^{-(2d+\varepsilon)}$$

- such that the generalized conditional expectation $E_X(a) := \text{Tr}_X(\gamma_X^* a \gamma_X)$ are completely positive, unital and KMS_{β} symmetric
- *L_X(a)* := *a* − *E_X(a)* is a bounded generators of a completely positive, unital, KMS_β symmetric semigroup
- a bounded Dirichlet form is given by

$$\mathcal{E}^{X}: L^{2}(\mathcal{M}, \omega) \to [0, +\infty) \qquad \mathcal{E}^{X}[i_{\omega}(\pi_{\omega}(a))] := \omega(\alpha_{-\frac{i\beta}{4}}^{\Phi}(a))\alpha_{-\frac{i\beta}{4}}^{\Phi}(L_{X}a))$$

• the quadratic form $\mathcal{E}: L^2(\mathcal{M}, \omega) \to [0, +\infty]$ $\mathcal{E}:=\sum_{j\in\mathbb{Z}^d} \mathcal{E}^{X+j}$ is densely defined, closable, Markovian and its closure is a Dirichlet form.

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Compact Quantum Groups							

A Compact Quantum Group $\mathbb{G} = (A, \Delta)$ is a unital C^{*}-algebra $A =: C(\mathbb{G})$ and

- a *coproduct* $\Delta : A \to A \otimes A$, a unital, *-homomorphism which is
- *coassociative* $(\Delta \otimes id_A) \circ \Delta = (id_A \otimes \Delta) \circ \Delta$ and satisfies
- cancelation rules $\overline{\text{Lin}}((1 \otimes A)\Delta(A)) = \overline{\text{Lin}}((A \otimes 1)\Delta(A)) = A \otimes A$.

A unitary corepresentation of \mathbb{G} is a unitary matrix $U = (u_{jk}) \in M_n(A)$ such that

•
$$\Delta(u_{jk}) = \sum_{p=1}^{n} u_{jp} \otimes u_{pk}$$
 $j, k = 1, \ldots, n$.

Theorem. (Woronowicz (1987))

Let $\{U^s : s \in \widehat{\mathbb{G}}\}\$ be a complete family of inequivalent irr. unitary corepr. of \mathbb{G} . Then the algebra of *polynomials*, defined by

$$\operatorname{Pol}(\mathbb{G}) := \operatorname{Span}\{u_{jk}^s; s \in \widehat{\mathbb{G}}, 1 \leq j, k \leq n_s\}$$

is a dense Hopf *-algebra with counit $\varepsilon(u_{jk}^s) := \delta_{jk}$ and antipode $S(u_{jk}^s) := (u_{kj}^s)^*$ satisfying $(m_A$ being the product in A)

 $(\varepsilon \otimes \mathrm{id})\Delta(a) = a \quad (\mathrm{id} \otimes \varepsilon)\Delta(a) = a \quad m_A(S \otimes \mathrm{id})\Delta(a) = \varepsilon(a)I = m_A(\mathrm{id} \otimes)\Delta(a) .$

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Convolution and Haar state							

• Convolution $\xi \star \xi' \in A^*$ of functionals $\xi, \xi' \in A^*$ is defined by

$$\xi \star \xi' := (\xi \otimes \xi') \circ \Delta;$$

• convolution $\xi \star a \in A$ of a functional $\xi \in A'$ and an element $a \in A$ is defined by

 $\xi \star a := (\mathrm{id} \otimes \xi)(\Delta a) \qquad a \star \xi := (\xi \otimes \mathrm{id})(\Delta a)$

Theorem. (Woronowicz (1987))

On a CQG $\mathbb{G} = (A, \Delta)$ there exists a unique (Haar) state $h \in A_+^*$ such that

$$h \star a = a \star h = h(a) \mathbf{1}_A \qquad a \in A$$
.

It is a $(\sigma, -1)$ -KMS state with respect to a suitable *-automorphisms group of A

$$\{\sigma_t : t \in \mathbb{R}\}$$
 $h(ab) = h(\sigma_{-i}(b)a)$ $a, b \in \mathcal{A}$.

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Notice that, in general, the Haar state is not a trace.

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Polar decomposition and unitary antipode							

Theorem. (Woronowicz (1987))

The antipode *S* is closable and its closure \overline{S} admits the polar decomposition:

$$\overline{S}=R\circ\tau_{\frac{i}{2}},$$

- $\tau_{\frac{1}{2}}$ generates a *-automorphisms group $\{\tau_t : t \in \mathbb{R}\}$ of the C^* -algebra A
- *R* is a linear, anti-multiplicative, norm preserving involution on *A* such that $\tau_t \circ R = R \circ \tau_t$ for all $t \in \mathbb{R}$, called *unitary antipode*.

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$SU_q(N)$							

Example: $SU_q(N)$

• The compact quantum group $SU_q(2) = (A, \Delta), 0 < q \le 1$, is given by the universal *C**-algebra *A* generated by the coefficients of the matrix

$$U = \left(egin{array}{cc} lpha & -q\gamma^* \ \gamma & lpha^* \end{array}
ight)$$

with relations on α and γ that ensuring unitarity $UU^* = U^*U = 1$

- comultiplication $\Delta(\alpha) := \alpha \otimes \alpha + \gamma \otimes \gamma, \quad \Delta(\gamma) := \gamma \otimes \alpha + \alpha^* \otimes \gamma$
- counit $\varepsilon(\alpha) = 1$ $\varepsilon(\gamma) = 0$
- antipode $S(\alpha) := \alpha^*$, $S(\gamma) := -q\gamma$, $S(u_{jk}^s) = (-q)^{(j-k)}u_{-k,-j}^s$
- *Haar state* $h(u_{jk}^s) = \delta_{s,0}$
- automorphisms group $\sigma_z(u_{jk}^s) = q^{2iz(j+k)}u_{jk}^s \quad z \in \mathbb{C}$
- *unitary antipode* $R(u_{jk}^s) = q^{k-j}(u_{kj}^s)^*$

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Let $\mathcal{A} = \text{Pol}(\mathbb{G})$ and (\mathcal{P}, Φ) a noncommutative probability space.

- **Random variable** on \mathcal{A} is a *-algebra homomorphism $j : \mathcal{A} \to \mathcal{P}$
- **distribution** of the random variable $j : \mathcal{A} \to \mathcal{P}$ is the state $\varphi_j = \Phi \circ j$
- **convolution** of the random variables $j_1, j_2 : \mathcal{A} \to \mathcal{P}$ is the random variable

$$j_1 \star j_2 = m_{\mathcal{P}} \circ (j_1 \otimes j_2) \circ \Delta$$

A Quantum Stochastic Process is a family of random variables $(j_{s,t})_{0 \le s \le t}$

- $j_{rs} \star j_{st} = j_{rt}$ for all $0 \le r \le s \le t \le T$ increment property and $j_{tt} = \varepsilon \mathbf{1}_{\mathcal{P}}$
- j_{st} converges to j_{ss} in distribution for $t \searrow s$ weak continuity.
- A QSP is called a Lévy Process if has
- independent increments, i.e. for disjoint intervals (*t_i*, *s_i*]

$$\Phi(j_{s_1t_1}(a_1)...j_{s_nt_n}(a_n)) = \Phi(j_{s_1t_1}(a_1))...\Phi(j_{s_nt_n}(a_n))$$

and $[j_{s_i,t_i}(a_1), j_{s_j,t_j}(a_2)] = 0$ for $i \neq j$,

• stationary increments, i.e. $\varphi_{st} = \Phi \circ j_{st}$ depends only on t - s,

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Theorem. (CFK 2011)

Lévy process $(j_{st})_{0 \le s \le t}$ on a A are in 1:1 correspondence with Markov semigroup (T_t) on A which are translation invariant

$$\Delta \circ T_t = (\mathrm{id} \otimes T_t) \circ \Delta \qquad t \ge 0.$$

"Proof". Distributions $\varphi_t := \varphi_{0,t} = \Phi \circ j_{0,t}$ form a continuous convolution semigroup of states on \mathcal{A} :

$$\varphi_0 = \varepsilon$$
 $\varphi_s \star \varphi_t = \varphi_{s+t}$ $\lim_{t \to 0} \varphi_t(b) = \varepsilon(b)$ $b \in \mathcal{A}$

whose generating functional $\varphi_t = \exp_{\star} tG$ is defined as $G = \frac{d}{dt}\varphi_t|_{t=0}$. A semigroup $T_t : \mathcal{A} \to \mathcal{A}$ is defined by the convolution

$$T_t = (\mathrm{id} \otimes \varphi_t) \circ \Delta = \varphi_t \star a, \quad t \ge 0$$

and its **infinitesimal generator** $L : A \to A$ results as the convolution operator associated to the generating functional

$$L(a) = (\mathrm{id} \otimes G) \circ \Delta(a) = G \star a$$
.

The semigroup extends to a translation invariant, Markov semigroup (T_i) on A and its generator is the closure of G. Moreover, has the relations

$$G = \varepsilon \circ L$$
, $\varphi_t = \varepsilon \circ T_t$ $t > 0$.

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KMS symmetric Lévy semigroups and spectrum								

Theorem. (C-Franz-Kula JFA 266 (2014))

Let $T_t = e^{-tL}$ be a Lévy semigroup on A with generating functional $G = \varepsilon \circ L$. The following properties are then equivalent

- the semigroup is KMS₋₁ symmetric with respect to the Haar state
- the generator is KMS₋₁ symmetric with respect to the Haar state
- the generating functional is invariant by the action of the unitary antipode R

 $G = G \circ R$ on the Hopf algebra $\mathcal{A} = \operatorname{Pol}(\mathbb{G})$.

Proposition. (C-Franz-Kula JFA 266 (2014))

 $L^{2}(A, h)$ decomposes as orthogonal sum of the finite dimensional subspaces

$$L^2(A,h) = \bigoplus_{s \in \widehat{\mathbb{G}}} E_s$$
 $E_s := \operatorname{Span} \{ u_{jk}^s \xi_h : j, k = 1, \cdots, n_s \}$ $s \in \widehat{\mathbb{G}}$.

L decomposes as a direct sum $L = \bigoplus_{s \in \widehat{\mathbb{G}}} L_s$ of its restrictions to the E_s subspaces. Its spectrum thus coincides with $\sigma(L) = \overline{\bigcup_{s \in \widehat{\mathbb{G}}} \sigma(L_s)}$.
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 Example: free orthogonal quantum group Orth

• The C*-algebra $C_u(O_N^+)$ is generated by $\{v_{jk} = v_{jk}^* : i, k = 1, \dots, N\}$ subject to

$$\sum_{l=1}^N v_{lj} v_{lk} = \delta_{jk} = \sum_{l=1}^N v_{jl} v_{kl} \qquad \Delta v_{jk} = \sum_{l=1}^N v_{lj} \otimes v_{lk}$$

• classes of irreducible, unitary corepresentations $O_N^+ \cong \mathbb{N}$

- the Haar *h* state is a trace, faithful on $Pol(O_N^+)$ but not on $C_u(O_N^+)$
- the Lévy semigroup e^{-tL} is constructed on the reduced C^{*}-algebra $C_r(O_N^+)$
- denote $U_s \in Pol[-N, N]$ the Chebyshev polynomial of the second kind

$$U_0(x) = 1, \quad U_1(x) = x, \quad U_s(x) = xU_{s-1}(x) - U_{s-1}(x), \qquad x \in [-N, N], \quad s \in \mathbb{N}$$

- generating functional $G(u_{jk}^{(s)}) := \delta_{jk} \frac{U'_s(N)}{U_s(N)}, \quad s \in \mathbb{N}, \quad j, k = 1, \cdots, U_s(N)$
- the generator has discrete spectrum, eigenvalues and multiplicities are given by

$$\lambda_s = rac{U_s'(N)}{U_s(N)}, \qquad m_s = \left(U_s(N)
ight)^2$$

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• spectral dimensions: $d_N = 3$ for N = 2, $d_N = +\infty$ for $N \ge 3$.

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Haagerup Approximation Property							

A second countable, locally compact group *G* has the Haagerup Approximation property HAP if there exists a sequence of normalized, positive definite functions $\varphi_n \in C_0(G)$, converging to the constant function 1 uniformly on compact subsets.

Equivalently, G has the HAP if there exists a proper, continuous, negative definite function on G.

By a result of U. Haagerup, the free groups \mathbb{F}_n have the HAP as their length functions are negative definite.

A long research (Connes-Jones, Choda, Jolissaint, Boca, Popa) culminated with various definitions of the HAP valid in general von Neumann algebras.

Let us consider the following one.

Definition. (Okayasu-Tomatsu 2014)

A von Neuman algebra \mathcal{M} has the HAP if there exists a standard form $(\mathcal{M}, \mathcal{H}, \mathcal{P}, \mathcal{J})$ and a sequence of contractive, completely positive operators $T_n : \mathcal{H} \to \mathcal{H}$ such that $\|\xi - T_n \xi\|_{\mathcal{H}} \to 0$ as $n \to +\infty$, for all $\xi \in \mathcal{H}$.

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HAP and I	Dirichlet forms						

Recently the above HAP has been found to be equivalent to others involving Markovian semigroups and Dirichlet forms.

Theorem. (Caspers-Skalski 2014)

The following properties are equivalent

- The von Neumann algebra \mathcal{M} has the HAP
- there exists a Markovian semigroup {*T_t* : *t* ≥ 0} w.r.t. a cyclic and separating vector ξ₀ ∈ *P*, such that *T_t* is compact for all *t* > 0
- there exists a Dirichlet form $(\mathcal{E}, \mathcal{F})$ w.r.t. a cyclic and separating vector $\xi_0 \in \mathcal{P}$, such that its spectrum is discrete.

As an application one can prove the following result.

Corollary. (Brannan 2012)

The von Neumann algebras $L^{\infty}(C_r(O_N^+), h)$ of the free orthogonal quantum groups O_N^+ in the cyclic representation of the Haar state h on $L^2(C_r(O_N^+), h)$, have Haagerup approximation property.

Proof. The result follows from the Caspers-Skalski equivalence and the construction of a Dirichlet form with discrete spectrum illustrated above.

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