

# Algebraic Renormalization Group and Pointlike Fields

Gerardo Morsella

Scuola Normale Superiore – Pisa

joint with H. Bostelmann, C. D'Antoni

Joint International Meeting UMI-DMV

Perugia, 18-22 June 2007

# Outline

- 1 Introduction
- 2 Background
  - Scaling Algebras
  - Pointlike Fields from Local Algebras
- 3 Scaling Limit of Pointlike Fields
  - Basic Idea
  - Phase Space and Scaling Limit
  - Renormalization Constants
- 4 Summary & Outlook

# Introduction

Conventional approach to the Renormalization Group:

- Pass from  $\phi$  to renormalized field  $\phi_\lambda(x) = Z_\lambda \phi(\lambda x)$ ;
- **Renormalization constants**  $Z_\lambda$  fixed by requiring e.g.  $\langle \Omega, \phi_\lambda(x) \phi_\lambda(y) \Omega \rangle \sim \text{const}$  as  $\lambda \rightarrow 0$ .

Problems:

- need to have detailed information on short-distance behaviour to calculate  $Z_\lambda$ ;
- fields do not have direct physical interpretation.

Algebraic approach [Buchholz-Verch '95]:

- $Z_\lambda$  not needed;
- based only on observables and model independent.

How does it compare to the conventional approach? How can we recover  $Z_\lambda$ ?

# Scaling Algebras

1/3

Data:

- $O \subset \mathbb{R}^4 \rightarrow \mathcal{A}(O) \subset B(\mathcal{H})$  net of observable algebras.
- $x \in \mathbb{R}^4 \rightarrow U(x)$  unitary representation of translations.
- $\Omega \in \mathcal{H}$  vacuum.

On  $C^*$ -algebra of bounded functions  $\lambda \in \mathbb{R}_+^{\times} \rightarrow \underline{A}_\lambda \in \mathcal{A}$

$$\alpha_x(\underline{A})_\lambda := \text{Ad } U(\lambda x)(\underline{A}_\lambda), \quad x \in \mathbb{R}^4,$$

Definition ([Buchholz-Verch '95])

**Local scaling algebra** of  $O$ :

$$\underline{\mathfrak{A}}(O) := \left\{ \underline{A} : \underline{A}_\lambda \in \mathcal{A}(\lambda O), \lim_{x \rightarrow 0} \|\alpha_x(\underline{A}) - \underline{A}\| = 0 \right\}$$

# Scaling Algebras

2/3

$\varphi$  locally normal state on  $\mathcal{A} \rightsquigarrow \underline{\varphi}_\lambda(\underline{A}) := \varphi(\underline{A}_\lambda)$  states on  $\underline{\mathcal{A}}$ ,

$SL^{\mathcal{A}}(\varphi) := \{\text{weak}^* \text{ limit points of } (\underline{\varphi}_\lambda)_{\lambda>0} \text{ for } \lambda \rightarrow 0\}$ .

## Theorem ([Buchholz-Verch '95])

- $SL^{\mathcal{A}}(\varphi)$  is independent of  $\varphi$ .
- $\underline{\omega}_0 \in SL^{\mathcal{A}}$  with GNS representation  $\pi_0$ . Then  $\mathcal{A}_0(O) := \pi_0(\underline{\mathcal{A}}(O))''$  is a covariant net in vacuum representation.

$O \rightarrow \mathcal{A}_0(O)$  is the **scaling limit net** of  $\mathcal{A}$ .

**Physical interpretation:**  $\mathcal{A}_0$  describes the short-distance (i.e. high-energy) behaviour of  $\mathcal{A}$ .

# Scaling Algebras

3/3

$\exists(\mathcal{F}, G)$  such that  $\mathcal{A} = \mathcal{F}^G$ . [Doplicher-Roberts '90]

**Field scaling algebra**  $\underline{\mathfrak{F}}$  and **scaling limit field net**  $\mathcal{F}_0$  defined in analogy to  $\underline{\mathfrak{A}}, \mathcal{A}_0$ .

$$\mathcal{A}_0 = \mathcal{F}_0^{G_0} \subset \mathcal{F}_0 \subseteq \mathcal{F}^0, \quad \mathcal{F}^0 \text{ DR net of } \mathcal{A}_0$$

$\mathcal{F}^0 \not\supseteq \mathcal{F}_0 \implies \mathcal{A}$  has **confined charges**.

E.g. in the **Schwinger model**  $\mathcal{F} = \mathcal{A} \implies \mathcal{F}_0 = \mathcal{A}_0 \subsetneq \mathcal{F}^0$ .

**Charge preservation**: identifies sectors of  $\mathcal{A}_0$  which are scaling limit of sectors of  $\mathcal{A}$  [D'Antoni-Verch-M. '03, D'Antoni-M. '06].

# Pointlike Fields from Local Algebras

1/3

**Basic idea** [Haag-Ojima '96]:

$$\Sigma_{E,r} = \{\sigma \upharpoonright \mathcal{A}(O_r) : \sigma \in P(E)B(\mathcal{H})_*P(E)\}$$

is compact for small  $r$

$\implies$  “finite” number of states describe short distance behaviour

$\implies$  basis  $(\phi_j)$  of  $\Sigma_{E,r}^*$  are pointlike fields.

# Pointlike Fields from Local Algebras

2/3

Quantitative version:

- $\Sigma = B(\mathcal{H})_*$ ,  $C^\infty(\Sigma) = \bigcap_{\ell > 0} R^\ell \Sigma R^\ell$ ,  $R = (1 + H)^{-1}$ ;
- $\|\sigma\|^{(\ell)} = \|R^{-\ell} \sigma R^{-\ell}\|$ ,  $\sigma \in C^\infty(\Sigma)$ ;
- $\Xi : \sigma \in C^\infty(\Sigma) \rightarrow \sigma \in \Sigma$ .

Definition ([Bostelmann '05])

$\mathcal{O} \rightarrow \mathcal{A}(\mathcal{O})$  satisfies the **asymptotic phase space condition** if  $\forall \gamma > 0$ ,  $\exists \ell > 0$ ,  $\psi : C^\infty(\Sigma) \rightarrow \Sigma$  of finite rank such that

$$\|\psi\|^{(\ell)} < \infty,$$

$$\|(\Xi - \psi) \upharpoonright \mathcal{A}(\mathcal{O}_r)\|^{(\ell)} = o(r^\gamma).$$



# Pointlike Fields from Local Algebras

3/3

rank  $\psi$  minimal,  $\psi = \sum_j \sigma_j \phi_j$ ,  $\sigma_j \in \Sigma$ ,  $\phi_j \in C^\infty(\Sigma)^*$ .

Define  $\Phi_\gamma := \text{span}\{\phi_j\}$ .  $\Phi_\gamma \subseteq \Phi_{\gamma'}$  if  $\gamma < \gamma'$ .

Theorem ([Bostelmann '05])

- $\Phi_\gamma$  independent of  $\psi$ ;
- $\cup_{\gamma>0} \Phi_\gamma = \{\phi : \exists \ell > 0, R^\ell \phi R^\ell \in \cap_{r>0} (R^\ell \mathcal{A}(O_r) R^\ell)^-\}$ .

$$\phi(f) = \int dx f(x) U(x) \phi U(x)^*, \quad \phi \in \Phi_\gamma,$$

Wightman field on  $C^\infty(H) = \cap_{\ell>0} R^\ell \mathcal{H}$ , and  $\phi(f) \eta \mathcal{A}(O)$   
[Fredenhagen-Hertel '81].

$\phi$  free:  $\Phi_0 = \mathbb{C}\mathbb{1}$ ,  $\Phi_1 = \text{span}\{\mathbb{1}, \phi\}$ ,  $\Phi_2 = \text{span}\{\Phi_1, \partial_\mu \phi, : \phi^2 : \}, \dots$

## Basic Idea

- Is the asymptotic phase space condition valid for  $\mathcal{A}_0$ ?
- Can we recover  $Z_\lambda$  such that  $\phi_0(x) = \lim_{\lambda \rightarrow 0} Z_\lambda \phi(\lambda x)$ ?

$\psi : C^\infty(\Sigma) \rightarrow \Sigma$  as above of rank 1:

$$\psi = \sigma\phi, \quad \sigma \in \Sigma, \phi \in \cup_{\gamma > 0} \Phi_\gamma.$$

Typically  $\|\sigma \upharpoonright \mathcal{A}(\lambda O)\| \rightarrow 0$  as  $\lambda \rightarrow 0$  (e.g. as  $O(\lambda)$  for free fields).

Let  $\underline{A} \in \underline{\mathcal{A}}(O)$ :  $\psi^*(\underline{A}_\lambda) = \sigma(\underline{A}_\lambda)\phi$  should be thought as a field at scale  $\lambda \implies$  we can choose  $Z_\lambda = \sigma(\underline{A}_\lambda) \sim \lambda$ .

**Message:** maps  $\psi$  are the good objects in the scaling limit.

# Phase Space and Scaling Limit

1/2

Scaling:  $r \rightarrow \lambda r$ ,  $E \rightarrow \lambda^{-1} E \implies$  phase space condition needs sharpening:

## Definition

$O \rightarrow \mathcal{A}(O)$  satisfies the **homogeneous asymptotic phase space condition** if  $\forall \gamma > 0$ ,  $\exists c, \ell > 0$  and  $\psi : C^\infty(\Sigma) \rightarrow \Sigma$  of finite rank such that

$$\begin{aligned} \|\psi \upharpoonright \Sigma_E, \mathcal{A}(O_r)\| &\leq c(1 + Er)^\ell, \\ \|(\Xi - \psi) \upharpoonright \Sigma_E, \mathcal{A}(O_r)\| &\leq c(Er)^\gamma. \end{aligned}$$

Satisfied by free fields.

Also by asymptotically free theories (logarithmic corrections)?

Note: HAPSC  $\implies$  APSC.

# Phase space and Scaling Limit

2/2

## Theorem

- $O \rightarrow \mathcal{A}(O)$  satisfies HAPSC  $\implies O \rightarrow \mathcal{A}_0(O)$  satisfies APSC;
- $\dim\Phi_{0,\gamma} \leq \dim\Phi_\gamma$ .

## Proof.

- Define  $\chi_0 : C^\infty(\Sigma_0) \rightarrow \mathfrak{A}^*$  by

$$\chi_0(\langle \pi_0(\underline{B})\Omega_0, (\cdot)\pi_0(\underline{B}')\Omega_0 \rangle)(\underline{A}) = \lim_\lambda \psi(\langle \underline{B}_\lambda\Omega, (\cdot)\underline{B}'_\lambda\Omega \rangle)(\underline{A}_\lambda),$$

where spectral support of  $\underline{B}, \underline{B}' = [0, E]$ ;

- $\text{rank}\chi_0 \leq \text{rank}\psi$ ;
- $\Phi_{0,\gamma} := (\ker\chi_0)^\perp$ ,  $\psi_0 := \Xi_0 \circ (\rho_{\Phi_{0,\gamma}})_*$ . □

# Renormalization Constants

- $\text{rank} \chi_0 < \infty \implies \Phi_{0,\gamma} = \text{img}(\chi_0^* \upharpoonright \underline{\mathfrak{A}})$ ;
- $\forall \phi_0 \in \Phi_{0,\gamma}, \exists \underline{A} \in \underline{\mathfrak{A}} : \phi_0 = \chi_0^*(\underline{A})$ ;
- Compute:

$$\begin{aligned}
 & \langle \pi_0(\underline{B})\Omega_0, \phi_0(x)\pi_0(\underline{B})\Omega_0 \rangle \\
 &= \chi_0(\langle \pi_0(\underline{\alpha}_{-x}\underline{B})\Omega_0, (\cdot)\pi_0(\underline{\alpha}_{-x}\underline{B})\Omega_0 \rangle)(\underline{A}) \\
 &= \lim_{\lambda} \psi(\langle (\underline{\alpha}_{-x}\underline{B})_{\lambda}\Omega, (\cdot)(\underline{\alpha}_{-x}\underline{B})_{\lambda}\Omega \rangle)(\underline{A}_{\lambda}) \\
 &= \lim_{\lambda} \sum_j \sigma_j(\underline{A}_{\lambda}) \langle \underline{B}_{\lambda}\Omega, \phi_j(\lambda x)\underline{B}_{\lambda}\Omega \rangle;
 \end{aligned}$$

- $Z_{j,\lambda} := \sigma_j(\underline{A}_{\lambda})$  **renormalization constants**.

This should be generalized to  $n$ -point functions.

## Summary & Outlook

In the algebraic approach to the Renormalization Group, the **field strength renormalization constants** are obtained in a **model-independent** way.

Outlook:

- **Coupling constants renormalization:** Can be replaced by scaling of OPE coefficients?
- **Asymptotic freedom:** Under which conditions the limit fields are free? Perhaps conditions on the 2-point functions?
- **Completeness:**  $\mathcal{A}$  generated by pointlike fields  $\implies \mathcal{A}_0$  generated by pointlike fields?