Coupled dynamics of two extended bodies

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Zielona Góra

- It's a very interesting city.
- Widely regarded as Europe's capital of art, architecture, music, shopping, sport and more.
- You should visit!



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Celestial bodies as spheres

- Celestial bodies are not spheres, but modelling them as such is convenient as the resulting equations of motion are integrable.
- However, certain phenomena, such as **gravitational torque**, are not described by this model.
- This is of importance to the description of double asteroids...
 - ... and furthermore for planning missions to asteroids.

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Extended bodies

- As a first approximation, take bodies to be rigid.
- One body may be well-approximated by a point-mass, depending on its geometry.
- Extended bodies may be well-approximated by **dumbbells**.
 - There are several variants of the dumbbell approach.
 - Only the very simplest variant is discussed in this talk.

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The dumbbell model: motivation

- Extended rigid bodies are modelled by two spheres separated by an inflexible massless rod.
- This model explains gravitational torque.
- (Weakness) Can only model **prolate** bodies, not **oblate** bodies, **unless** complex masses and complex rod lengths are used.

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Point-mass and solid dumbbell

- One way in which dumbbells can be classified is by whether the two spheres with which the dumbbell is composed have finite or zero radius.
- The zero radius case sometimes results in fewer equations of motion; however, it must be treated differently as the inertia tensor of such a body is singular in \mathbb{R}^3 . This is the case that this discussion concerns.

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State of literature

- Equations of motion are derived.
- Relative equilibria are identified. However...
- Bifurcation analysis of libration points is convoluted and incomplete.
- Stability analysis of libration points is similarly deficient.
- Invariant manifolds are not determined.
- Question of integrability remains unanswered.

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Our contribution

- A novel global analysis of the point-mass dumbbell + point-mass problem, including:
 - An exhaustive analysis of the libration points, including their stability and any bifurcations that arise.
 - Proofs of the (non-)integrability of both the general point-mass dumbbell problem and the special case where the dumbbell is symmetric.
- Eventually, a similarly complete analysis for the more complex(!) dumbbell systems.

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Equations of motion

$$\dot{R}_1 = P_1 - I^{-1} K_1 R_3, \tag{1}$$

$$\dot{R}_3 = P_3 + I^{-1} K_1 R_1, \tag{2}$$

$$\dot{P}_{1} = \frac{P_{2}^{2}}{R_{1}} - I^{-1} \left(\frac{K_{2}R_{3}P_{2}}{R_{1}} + K_{1}P_{3} \right) - \frac{\mu(1+\mu)^{2}R_{1}}{\left\| (1+\mu)\mathbf{R} - \begin{pmatrix} 0 \\ 0 \end{pmatrix} \right\|^{3}} - \frac{(1+\mu)^{2}R_{1}}{\left\| (1+\mu)\mathbf{R} + \mu \begin{pmatrix} 0 \\ 0 \end{pmatrix} \right\|^{3}}, \quad (3)$$

$$\dot{P}_2 = -\frac{P_1 P_2}{R_1} + I^{-1} K_2 \frac{R_3 P_1 - R_1 P_3}{R_1},$$
(4)

$$\dot{P}_{3} = I^{-1} \left(K_{1} P_{1} + K_{2} P_{2} \right) - \frac{\mu (1+\mu)(-1+R_{3}+R_{3}\mu)}{\left\| (1+\mu)\mathbf{R} - \begin{pmatrix} 0\\ 0\\ 1 \end{pmatrix} \right\|^{3}} - \frac{(1+\mu)(R_{3}+\mu+R_{3}\mu)}{\left\| (1+\mu)\mathbf{R} + \mu \begin{pmatrix} 0\\ 0\\ 1 \end{pmatrix} \right\|^{3}},$$
(5)

$$\dot{K}_{1} = \frac{-I^{-1}K_{2}^{2}R_{3} + K_{2}P_{2}}{R_{1}} + \frac{\mu(1+\mu)R_{1}}{\left\| (1+\mu)\mathbf{R} - \begin{pmatrix} 0\\ 0\\ 1 \end{pmatrix} \right\|^{3}} - \frac{\mu(1+\mu)R_{1}}{\left\| (1+\mu)\mathbf{R} + \mu \begin{pmatrix} 0\\ 0\\ 1 \end{pmatrix} \right\|^{3}},$$
(6)
$$\dot{K}_{2} = I^{-1}K_{1}\frac{K_{2}R_{3} - IP_{2}}{R_{1}}.$$
(7)

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QUESTIONS?

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