F. Gavarini

"The global quantum duality principle: a survey through examples"

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INTRODUCTION

"Dualitas dualitatum et omnia dualitas"

N. Barbecue, "Scholia"

The most general notion of "symmetry" in mathematics is encoded in the notion of Hopf algebra. Among Hopf algebras over a field, the commutative and the cocommutative ones encode "geometrical" symmetries, in that they correspond, under some technical conditions, to algebraic groups and to (restricted, if the ground field has positive characteristic) Lie algebras respectively: in the first case the Hopf algebra is the algebra F[G] of regular functions over an algebraic group G, whereas in the second case it is the (restricted) universal enveloping algebra $U(\mathfrak{g})$ ($\mathfrak{u}(\mathfrak{g})$ in the restricted case) of a (restricted) Lie algebra g. A popular generalization of these two types of "geometrical symmetry" is given by quantum groups: roughly, these are Hopf algebras H depending on a parameter \hbar such that setting $\hbar = 0$ the Hopf algebra one gets is either of the type F[G] — hence H is a quantized function algebra, in short QFA — or of the type $U(\mathfrak{g})$ or $\mathbf{u}(\mathfrak{g})$ (according to the characteristic of the ground field) — hence H is a quantized universal enveloping algebra, in short QrUEA. When a QFA exists whose specialization (i.e. its "value" at $\hbar = 0$) is F[G], the algebraic group G inherits from this "quantization" a Poisson bracket, which makes it into a Poisson (algebraic) group; similarly, if a QrUEA exists whose specialization is $U(\mathfrak{g})$ or $\mathbf{u}(\mathfrak{g})$, the (restricted) Lie algebra \mathfrak{g} inherits a Lie cobracket which makes it into a Lie bialgebra. Then by Poisson group theory one has Poisson groups G^* dual to G and a Lie bialgebra \mathfrak{g}^* dual to \mathfrak{g} , so other geometrical symmetries are related to the initial ones.

The dependence of a Hopf algebra on \hbar can be described as saying that it is defined over a ring R and $\hbar \in R$: so one is lead to dwell upon the category \mathcal{HA} of Hopf R-algebras (maybe with some further conditions), and then raises three basic questions:

- (1) How can we produce quantum groups?

- (2) How can we characterize quantum groups (of either kind) within \mathcal{HA} ?

— (3) What kind of relationship, if any, does exist between quantum groups over mutually dual Poisson groups, or mutually dual Lie bialgebras?

A first answer to question (1) and (3) together is given by the so-called "quantum duality principle", which is known in literature with at least two formulations. One claims that quantum function algebras associated to dual Poisson groups can be taken to be dual — in the Hopf sense — to each other; and similarly for quantum enveloping algebras (cf. [FRT1] and [Se]). The second one, formulated by Drinfeld in local terms (i.e., using

formal groups — rather than algebraic groups — and complete topological Hopf algebras; cf. [Dr], §7, and see [Ga4] for a proof), provides a recipe to get, out of a QFA over G, a QrUEA over \mathfrak{g}^* , and, conversely, to get, out of a QrUEA over \mathfrak{g} , a QFA over G^* . More precisely, Drinfeld defines two functors, inverse to each other, from the category of quantized universal enveloping algebras (in his sense) to the category of quantum formal series Hopf algebras (Drinfeld's analogue of QFAs) and viceversa, such that $U_{\hbar}(\mathfrak{g}) \mapsto$ $F_{\hbar}[[G^*]]$ and $F_{\hbar}[[G]] \mapsto U_{\hbar}(\mathfrak{g}^*)$ (in his notation, where the subscript \hbar stands as a reminder for "quantized" and the double square brackets stand for "formal series Hopf algebra").

In this paper we provide a global version of the quantum duality principle which gives a complete answer to questions (1) through (3). The key is to push as far as possible Drinfeld's original idea so to apply it to the category \mathcal{HA} of all Hopf algebras which are torsion-free modules over some 1-dimensional domain (in short, 1dD), say R, and to do it for each non-zero prime element \hbar in R.

To be precise, we extend Drinfeld's recipe so to define functors from \mathcal{HA} to itself. We show that the image of these "generalized" Drinfeld's functors is contained in a category of quantum groups — one gives QFAs, the other QrUEAs — so we answer question (1). Then, in the zero characteristic case, when restricted to quantum groups these functors yield equivalences inverse to each other. Moreover these equivalences exchange the types of quantum group — switching QFA with QrUEA — and the underlying Poisson symmetries — interchanging G or \mathfrak{g} with G^* or \mathfrak{g}^* , thus solving (3). Other details enter the picture to show that these functors endow \mathcal{HA} with sort of a (inner) "Galois' correspondence", in which QFAs on one side and QrUEAs on the other side are the subcategories (in \mathcal{HA}) of "fixed points" for the composition of both Drinfeld's functors (in the suitable order): in particular, this answers question (2). It is worth stressing that, since our "Drinfeld's functors" are defined for each non-trivial point (\hbar) of Spec(R), for any such (\hbar) and for any H in \mathcal{HA} they yield two quantum groups, namely a QFA and a QrUEA, w.r.t. \hbar itself. Thus we have a method to get, out of any single $H \in \mathcal{HA}$, several quantum groups.

Therefore the "global" in the title is meant in several respects: geometrically, we consider global objects (namely Poisson groups rather than Poisson formal groups, as in Drinfeld's approach); algebraically we consider quantum groups over any 1dD R, so there may be several different "semiclassical limits" (=specialization) to consider, one for each non-trivial point of the spectrum of R (whereas in Drinfeld's context $R = \&[[\hbar]]$ so one can specialise only at $\hbar = 0$); more generally, our recipe applies to any Hopf algebra, i.e. not only to quantum groups; finally, the most part of our results are characteristic-free, i.e. they hold not only in zero characteristic (as in Drinfeld's case) but also in positive characteristic.

As a further outcome, this "global version" of the quantum duality principle opens the way to formulate a "quantum duality principle for subgroups and homogeneous spaces": this is the subject of a separate paper (cf. [CG]).

A key, long-ranging application of our "global quantum duality principle" is the following. Take as R the polynomial ring $R = \Bbbk[\hbar]$, where \Bbbk is a field: then for any Hopf algebra over \Bbbk we have that $H[\hbar] := R \otimes_{\Bbbk} H$ is a torsion-free Hopf R-algebra, hence we can apply Drinfeld's functors to it. The outcome of this procedure is the "crystal duality principle", whose statement strictly resembles that of the quantum duality principle: now Hopf \Bbbk -algebras are looked at instead of torsionless Hopf R-algebras, and quantum groups are replaced by Hopf algebras with canonical filtrations such that the associated graded Hopf algebra is either commutative or cocommutative. Correspondingly, we have a method to associate to H a Poisson group G and a Lie bialgebra \mathfrak{k} such that G is an affine space (as an algebraic variety) and \mathfrak{k} is graded (as a Lie bialgebra); in both cases, the "geometrical" Hopf algebra can be attained — roughly — through a continuous 1-parameter deformation process. This result can also be formulated in purely classical (i.e. "non-quantum") terms and proved by purely classical means. However, the approach via the global quantum duality principle also yields further possibilities to deform H into other Hopf algebras of geometrical type, which is out of reach of any classical approach.

The purpose of these notes is to illustrate the global quantum duality principle in some detail through some relevant examples: some quantum groups — the standard quantization of the Kostant-Kirillov structure on a Lie algebra ($\S4$), the quantum semisimple groups ($\S5$), the three dimensional quantum Euclidean group ($\S6$), the quantum Heisenberg group, and a special, far-reaching application, the "crystal duality principle" ($\S3$). All details and technicalities which are skipped in the present paper can be found in [Ga5].

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